

## BOUNDED EIGENFUNCTIONS AND ABSOLUTELY CONTINUOUS SPECTRA FOR ONE-DIMENSIONAL SCHRÖDINGER OPERATORS

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(Communicated by Palle E. T. Jorgensen)

ABSTRACT. We provide a short proof of that case of the Gilbert-Pearson theorem that is most often used: That all eigenfunctions bounded implies purely a.c. spectrum. Two appendices illuminate Weidmann's result that potentials of bounded variation have strictly a.c. spectrum on a half-axis.

### §1. INTRODUCTION AND REDUCTION TO $m$ -FUNCTIONS

In this note, I want to consider Schrödinger operators and Jacobi matrices on a half-line. Specifically, we'll consider the operator  $h$  on  $\ell^2(\mathbb{Z}_+)$  (with  $\mathbb{Z}_+ = \{1, 2, \dots\}$ ) given by

$$(1.1a) \quad (hu)(n) = u(n+1) + u(n-1) + v(n)u(n),$$

$$(1.1b) \quad u(0) = 0,$$

and the self-adjoint operator on  $L^2(0, \infty)$

$$(1.2a) \quad (Hu)(x) = -u''(x) + V(x)u(x),$$

$$(1.2b) \quad u(0) = 0,$$

where we suppose

$$(1.3) \quad \Gamma(V) \equiv \sup_x \left( \int_{x-1}^{x+1} |V(y)|^2 \right) < \infty.$$

For any  $E \in \mathbb{C}$ , define two solutions  $u_1, u_2$  of the formal difference (resp. differential) equation  $hu = Eu$  (resp.  $Hu = Eu$ ) with boundary conditions:

$$u_1(0, E) = 0, \quad u_1(1, E) = 1,$$

$$u_2(0, E) = 1, \quad u_2(1, E) = 0,$$

in the discrete case and

$$u_1(0, E) = 0, \quad u_1'(0, E) = 1,$$

$$u_2(0, E) = 1, \quad u_2'(0, E) = 0,$$

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Received by the editors April 3, 1995 and, in revised form, April 24, 1995.

1991 *Mathematics Subject Classification*. Primary 34L40.

This material is based upon work supported by the National Science Foundation under Grant No. DMS-9401491. The Government has certain rights in this material.

in the continuous case.

Let  $S = \{E \in \mathbb{R} \mid u_1 \text{ and } u_2 \text{ are bounded on } [0, \infty)\}$ . Then our purpose here is to prove

**Theorem 1.** *On  $S$ , the spectral measure  $\rho$  for  $h$  (resp.  $H$ ) is purely absolutely continuous in the sense that*

- (i)  $\rho_{ac}(T) > 0$  for any  $T \subset S$  with  $|T| > 0$  (where  $|\cdot| = \text{Lebesgue measure}$ ).
- (ii)  $\rho_{sing}(S) = 0$ .

This theorem is not new. In [9], [8], [11], [14], Gilbert, Kahn, and Pearson proved a complete characterization of the essential support of  $\rho_{ac}$  in terms of mutually subordinate solutions. Their approach has the advantage of not requiring (1.3). Behncke [2] and Stolz [17] have noted that  $V$  uniformly  $L^1_{loc}$  with bounded eigenfunctions allows one to use the Gilbert-Pearson theory. Virtually all applications of [9], [11] use the weaker Theorem 1. There seems to be some point in the short proof I'll present here which avoids some of their tricky calculations and which makes the result transparent. In addition, we'll obtain explicit bounds on  $m$ -functions.

I should mention earlier work of Carmona [4] (which is weaker than Theorem 1) and related work of Briet-Mourre [3].

As with Gilbert-Pearson, our proof uses the theory of Weyl  $m$ -functions. For  $E \in \mathbb{C}_+ = \{z \mid \text{Im } z > 0\}$ , we can find a unique solution  $u_+(n, E)$  (resp.  $u_+(x, E)$ ) of (1.1a)/(1.2a) with  $u_+ \in \ell^2$  (resp.  $L^2$ ) at infinity, normalized by

$$(1.4) \quad u_+(0, E) = 1.$$

Then one defines the  $m$  function by

$$(1.5) \quad m_+(E) = u_+(1, E)$$

in the discrete case and

$$(1.6) \quad m_+(E) = u'_+(0, E)$$

in the continuous case.

By looking at the Wronskian of  $u_+$  and  $\bar{u}_+$ , one gets the well-known formula:

$$(1.7) \quad \text{Im } m_+(E) = \text{Im } E \sum_{n=1}^{\infty} |u_+(n, E)|^2$$

in the discrete case and

$$(1.8) \quad \text{Im } m_+(E) = \text{Im } E \int_0^{\infty} |u_+(x, E)|^2 dx$$

in the continuous case.

It is known (see [5], [16]) that

$$(1.9) \quad d\rho(E) = \frac{1}{\pi} \lim_{\epsilon \downarrow 0} \text{Im } m_+(E + i\epsilon) dE.$$

It follows [1], [7] by the de la Vallée-Poussin theorem that

$$\rho_{sing} \text{ is supported on } \left\{ E \mid \lim_{\epsilon \downarrow 0} \text{Im } m_+(E + i\epsilon) = \infty \right\}$$

and

$$d\rho_{ac}(E) = \frac{1}{\pi} \operatorname{Im} m_+(E + i0) dE.$$

Thus, Theorem 1 is an immediate consequence of

**Theorem 2.** *If  $E \in S$ , then*

$$(1.10) \quad (i) \quad \underline{\lim} \operatorname{Im} m_+(E + i0) > 0,$$

$$(1.11) \quad (ii) \quad \overline{\lim} |m_+(E + i0)| < \infty.$$

*Remark.* While the results are stated for the half-line with Dirichlet boundary conditions, Theorem 2 immediately implies the result for any fixed boundary condition and for the whole line. For it is known [1], [16] that the essential support  $d\rho_{ac,\theta}$  for  $\theta$  boundary conditions (given by  $\sin(\theta)u'(0) + \cos(\theta)u(0) = 0$ ) is  $\theta$  independent and that  $d\rho_{sing,\theta}$  is supported on the set where  $m_+(E + i0) = -\cot(\theta)$ , which cannot happen if (1.10)/(1.11) holds. For the whole line, we can define  $S$  via the right half-line condition from which (1.10)/(1.11) and the formulas (for the continuous case; the discrete case is similar)

$$d\rho_1(E) = -\lim_{\epsilon \downarrow 0} \frac{1}{\pi} \operatorname{Im} \left( \frac{1}{m_+(E + i\epsilon) + m_-(E + i\epsilon)} \right) dE,$$

$$d\rho_2(E) = \lim_{\epsilon \downarrow 0} \frac{1}{\pi} \operatorname{Im} \left( \frac{1}{m_+(E + i0)^{-1} + m_-(E + i0)^{-1}} \right) dE$$

imply  $\rho_{i,sing}(S) = 0$ .

It is a pleasure to thank F. Gesztesy, A. Kiselev, and G. Stolz for useful discussions.

### §2. THE JACOBI MATRIX CASE

In this section, we'll prove Theorem 2 in the discrete case. Define the fundamental or transfer matrix by

$$T(E, n, 0) = \begin{pmatrix} u_1(n + 1, E) & u_2(n + 1, E) \\ u_1(n, E) & u_2(n, E) \end{pmatrix}$$

and then

$$(2.1) \quad T(E, n, m) = T(E, n, 0)T(E, m, 0)^{-1}.$$

$T$  is defined so that if  $u$  obeys  $hu = Eu$ , then  $\Phi(n) = \begin{pmatrix} u(n+1) \\ u(n) \end{pmatrix}$  obeys

$$\Phi(n) = T(E, n, m)\Phi(m).$$

Constancy of the Wronskian implies  $\det T = 1$  so  $\|T^{-1}\| = \|T\|$  and thus by (2.1)

$$(2.2) \quad C(E) = \sup_{n,m} \|T(E, n, m)\| \leq \sup_n \|T(E, n, 0)\|^2$$

is finite if and only if  $E \in S$ . We'll prove Theorem 2 in the following explicit form:

**Theorem 2J.** *If  $E \in S$ , then*

$$(2.3) \quad \underline{\lim} \operatorname{Im} m_+(E + i\epsilon) \geq \frac{1}{4} C^{-3},$$

$$(2.4) \quad \overline{\lim} |m_+(E + i\epsilon)| \leq 4C^3,$$

where  $C(E)$  is given by (2.2).

*Proof.* Let

$$A(E, n) \equiv \begin{pmatrix} E - V(n) & -1 \\ 1 & 0 \end{pmatrix}$$

so  $T(E, n, 0) = A(E, n)T(E, n - 1, 0)$ . It follows (as a telescoping sum) that

$$T(E + i\epsilon, n, 0) = T(E, n, 0) + \sum_{j=0}^{n-1} (i\epsilon)T(E, n, j + 1) \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} T(E + i\epsilon, j, 0)$$

so by iteration, we get

$$(2.5) \quad \|T(E + i\epsilon, n, 0)\| \leq \sum_{k=0}^n \binom{n}{k} C^{k+1} \epsilon^k = C(1 + C\epsilon)^n \leq Ce^{\epsilon Cn}.$$

By  $\|T^{-1}\| = \|T\|$ , we see that

$$\left\| \begin{pmatrix} u_+(E + i\epsilon, n + 1) \\ u_+(E + i\epsilon, n) \end{pmatrix} \right\| \geq C^{-1} e^{-\epsilon Cn} (|m_+(E + i\epsilon)|^2 + 1)^{1/2}$$

since  $\begin{pmatrix} u_+(E+i\epsilon, 1) \\ u_+(E+i\epsilon, 0) \end{pmatrix} = \begin{pmatrix} m_+(E+i\epsilon) \\ 1 \end{pmatrix}$ .

Squaring and summing over  $n = 1, 3, \dots$  we see that

$$\sum_{n=1}^{\infty} |u_+(E + i\epsilon, n)|^2 \geq C^{-2} e^{-2\epsilon C} (1 - e^{-4\epsilon C})^{-1} (|m_+(E + i\epsilon)|^2 + 1).$$

Thus by (1.7)

$$\text{Im } m_+(E + i\epsilon) \geq \frac{1}{4} C^{-3} e^{-2\epsilon C} (4\epsilon C) (1 - e^{-4\epsilon C})^{-1} [1 + |m_+(E + i\epsilon)|^2]$$

or

$$(2.6) \quad \underline{\lim} \left[ \text{Im } m_+(E + i\epsilon) / [1 + |m_+(E + i\epsilon)|^2] \right] \geq \frac{1}{4} C^{-3}.$$

Noting that  $(1 + |m_+|^2)^{-1} \leq 1$ , we see that (2.6) immediately implies (2.3). And since  $(1 + |m_+|^2) / \text{Im } m_+ \geq |m_+|$ , it also implies (2.4).  $\square$

With only minor changes, the theorem extends to the general Jacobi matrix (tridiagonal self-adjoint) matrix:

$$(2.7) \quad (hu)(n) = a_{n+1}u(n + 1) + a_nu(n - 1) + b_nu(n)$$

so long as there is  $\alpha$  finite with

$$(2.8) \quad \alpha^{-1} < |a_n| < \alpha$$

for all  $n$ . If  $d\rho$  is the spectral measure for  $u(n) = \delta_{1n}$ , then

$$\int \frac{d\rho(E)}{z - E} = m_+(z),$$

where  $m_+(z)$  is defined to be  $a_1^{-1}u_+(1)$  (if  $u_+$  is normalized by  $u_+(0) = 1$ ). (1.7) becomes

$$\text{Im } m_+(E) = a_1^{-2} (\text{Im } E) \sum_{n=1}^{\infty} |u_+(n, E)|^2.$$

It is no longer true that  $\|T(E, n, 0)^{-1}\| = \|T(E, n, 0)\|$  since  $\det(T(E, n, 0))$  may not be 1. Rather  $\det(T(E, n, 0)) = \frac{\alpha_1}{a_{n+1}}$  so using (2.8), (2.2) becomes  $C(E) \leq \alpha^2 \sup_n \|T(E, n, 0)\|^2$ . (2.5) becomes

$$\|T(E + i\epsilon, n, 0)\| \leq C(1 + C\alpha\epsilon)^n \leq Ce^{\epsilon Cn\alpha}$$

and (2.6) becomes

$$\underline{\lim} [\operatorname{Im} m_+(E + i\epsilon)] / [1 + a_1^2 |m_+(E + i\epsilon)|^2] \geq \frac{1}{4} a_1^{-2} C^{-3} \alpha^{-1}.$$

### §3. THE SCHRÖDINGER CASE

To carry the proof through from the discrete case, we must use (1.3) to bound  $u'$  locally by  $u$ . This is a standard Sobolev-type estimate; we haven't tried to optimize constants.

**Lemma 3.1.** *If  $u$  obeys  $-u'' + Vu = Eu$ , then*

$$(3.1) \quad |u'(x)|^2 \leq \left[ 4 + \frac{3}{4} \Gamma(|V - E|) \right] \int_{x-1}^{x+1} |u(y)|^2 dy$$

where  $\Gamma$  is given by (1.3).

*Proof.* By Taylor's theorem with remainder,

$$f'(0) = \frac{1}{2} \frac{f(x) - f(-x)}{x} - \frac{1}{2x} \int_0^x (x - y)[f''(y) + f''(-y)] dy.$$

Integrate this from  $\frac{1}{2}$  to 1 to get

$$|f'(0)| \leq \int_{-1}^1 |f(x)| dx + \frac{3}{8} \int_{-1}^1 |f''(x)| dx.$$

Let  $f(y) = u(y + x)$  and use  $u'' = (V - E)u$  and the Schwarz inequality to get (3.1). □

By (3.1), if  $E \in S$ ,  $u'$  is also bounded and thus the transfer matrix  $T(E, x, y)$  defined by

$$T(E, x, y) \begin{pmatrix} u'(y) \\ u(y) \end{pmatrix} = \begin{pmatrix} u'(x) \\ u(x) \end{pmatrix}$$

is bounded. Let

$$C(E) \equiv \sup_{x,y} \|T(E, x, y)\|.$$

**Theorem 2S.** *Let  $E \in S$  and define  $A(E) = \frac{1}{2} C(E)^{-3} / (9 + \frac{3}{2} \Gamma(|E - V|))$ . Then*

$$\underline{\lim} \operatorname{Im} m_+(E + i\epsilon) \geq A, \\ \overline{\lim} |m_+(E + i\epsilon)| \leq A^{-1}.$$

*Proof.* By mimicking the proof of (2.5), using integrals in place of sums,

$$(3.2) \quad \|T(E + i\epsilon, x, 0)\| \leq C e^{\epsilon C|x|}.$$

By (3.1)

$$\int_1^\infty |u'(x)|^2 dx \leq \left[8 + \frac{3}{2} \Gamma(|V - E - i\epsilon|)\right] \int_0^\infty |u(y)|^2 dy$$

so if  $\beta = 1/(9 + \frac{3}{2}\Gamma)$ , then

$$\begin{aligned} \int_0^\infty |u(y)|^2 dy &\geq \beta \int_1^\infty [|u(x)|^2 + |u'(x)|^2] dx \\ &\geq C^{-2} \beta (1 + |m_+|^2) \int_1^\infty e^{-2\epsilon Cx} dx \end{aligned}$$

so by (1.8),

$$\text{Im } m_+ \geq \frac{1}{2} C^{-3} \beta e^{-\epsilon C} (1 + |m_+|^2)$$

and the result follows as in the discrete case. □

APPENDIX 1: A DISCRETE VERSION OF WEIDMANN'S THEOREM

One of the more interesting applications of Theorem 2 is the result of Weidmann [18], [19], [20] that if  $V = V_1 + V_2$ , where  $V_1 \in L^1$  and  $V_2$  is of bounded variation with  $V_2(x) \rightarrow 0$  at infinity, then  $-\frac{d^2}{dx^2} + V(x)$  has purely a.c. spectrum on  $(0, \infty)$ . A key to his argument is a proof that for any  $E > 0$ , solutions are bounded. He does this by noting one can suppose  $V_2$  is  $C^1$  with  $V_2' \in L^1$  (by adjusting the breakup) and that if  $K(x) = (u')^2 + (E - V_2)u^2$ , then  $K'(x) = 2V_1u'u - 2V_2'u^2 \leq C(|V_1| + |V_2'|)K(x)$  for  $x$  large. Here we'll prove a result of Máté and Nevai [13] using a discrete analog of Weidmann's approach:

**Theorem A.1.** *Let  $v_n$  be a sequence on  $\{1, 2, \dots\}$  so that  $v_n \rightarrow 0$  and*

$$(A.1) \quad \sum_{n=1}^\infty |v_{n+1} - v_n| < \infty.$$

*Then, the operator  $h$  of (1.1) has purely absolutely continuous spectrum on  $(-2, 2)$ .*

*Remarks.* 1. (A.1) implies  $\lim v_n$  exists so by adding a constant, it is no loss to suppose  $v_n \rightarrow 0$ .

2. If  $v_n \in \ell^1$ , then (A.1) holds so we don't need to consider sums as Weidmann does in the continuous case.

*Proof.* Given a solution of  $hu = Eu$ , let

$$K_n = u_{n+1}^2 + u_n^2 + (v_n - E)u_n u_{n+1}.$$

Then

$$(K_{n+1} - K_n) = (u_{n+2} - u_n)(u_{n+2} + u_n + (v_{n+1} - E)u_{n+1}) + (v_n - v_{n+1})u_n u_{n+1}$$

so

$$(A.2) \quad |K_{n+1} - K_n| \leq |v_n - v_{n+1}| |u_n u_{n+1}|.$$

Suppose now  $E \in (-2, 2)$ . Then for  $n \geq$  some  $N_0$ ,  $2 - |v_n - E| \geq \delta > 0$ . For such  $n$ ,

$$K_n \geq \frac{\delta}{2} (u_{n+1}^2 + u_n^2) + \left(1 - \frac{\delta}{2}\right) (|u_{n+1}| - |u_n|)^2 \geq \frac{\delta}{2} (u_{n+1}^2 + u_n^2)$$

so (A.2) becomes

$$K_{n+1} \leq \left(1 + \frac{2}{\delta} |v_n - v_{n+1}|\right) K_n$$

and for all  $n \geq N_0$ :

$$K_n \leq \prod_{m=N_0}^{\infty} \left(1 + \frac{2}{\delta} |v_m - v_{m+1}|\right) K_{N_0}.$$

The product is convergent by (A.1). □

By using the remark at the end of Section 1, Theorem A.1 extends to the operator (2.7) so long as (2.8) holds and

$$\begin{aligned} b_n &\rightarrow 0, & \sum_{n=1}^{\infty} |b_{n+1} - b_n| &< \infty, \\ a_n &\rightarrow 1, & \sum_{n=1}^{\infty} |a_{n+1} - a_n| &< \infty. \end{aligned}$$

We merely define  $K_n$  by

$$K_n = a_{n+1} u_{n+1}^2 + a_n u_n^2 + (b_n - E) e_n u_{n+1}.$$

This is related to results of [6].

### APPENDIX 2: EIGENFUNCTIONS FOR WEIDMANN'S THEOREM

We want to further elucidate Weidmann's theorem by showing how to actually find the asymptotics of the eigenfunctions. We'll suppose  $V(x) = V_1(x) + V_2(x)$  with  $V_1 \in L^1$  and  $V_2$  a  $C^1$  function with  $V_2' \in L^1$  and  $V_2 \rightarrow 0$  at infinity. We claim:

**Theorem B.1.** *Fix  $E = k^2 > 0$  with  $k > 0$ . Then every solution of*

$$\left(-\frac{d^2}{dx^2} + V(x)\right) u = Eu$$

*is bounded; indeed, there exist  $a, b$  so that*

$$\begin{aligned} |u(x) - au_+(x) - bu(x)| &\rightarrow 0, \\ |u'(x) - iaku_+(x) + ibku_-(x)| &\rightarrow 0, \end{aligned}$$

*where*

$$(B.1) \quad u_{\pm}(x) = \exp\left(\pm i \int_{x_0}^x \sqrt{k^2 - V_2(x)} dx\right)$$

*and  $x_0$  is chosen so large that  $V_2(x) < k^2$  for  $x > x_0$ .*

*Remarks.* 1. Since  $(k^2 - V_2(x))^{-1/4} \rightarrow k^{-1/2}$ , we could use the WKB form instead of (B.1), but the form (B.1) is what enters naturally.

2. This theorem and proof can be regarded as specializations of arguments in Hinton-Shaw [10].

*Proof.* Define  $u_{\pm}$  by (B.1). Note that  $u_{\pm}$  are  $C^2$  and

$$(B.2a) \quad -u''_{\pm} + (V(x) - E)u_{\pm} = F_{\pm}u_{\pm},$$

where

$$(B.2b) \quad F_{\pm}(x) = V_1(x) \pm \frac{i}{2} V_2'(x)(k^2 - V_2(x))^{-1/2}$$

is in  $L^1$  near infinity.

Let  $W(x)$  be the Wronskian of  $u_+$  and  $u_-$ . Clearly,  $W(x) = 2ik + o(1)$ . Define  $a(x), b(x)$  by the equations (variation of parameters)

$$\begin{aligned} u(x) &= a(x)u_+(x) + b(x)u_-(x), \\ u'(x) &= a(x)u'_+(x) + b(x)u'_-(x). \end{aligned}$$

A straightforward and standard calculation (see prob. 98 on pg. 395 of [15]) shows that  $a, b$  obey the equations

$$\begin{pmatrix} a(x) \\ b(x) \end{pmatrix}' = M(x) \begin{pmatrix} a(x) \\ b(x) \end{pmatrix}$$

where

$$M(x) = W(x)^{-1} \begin{pmatrix} -F_+ & -u^2 F_- \\ F_+ u_+^2 & F_- \end{pmatrix}.$$

Since this is in  $L^1$ , standard arguments show that  $\lim_{k \rightarrow \infty} \begin{pmatrix} a(x) \\ b(x) \end{pmatrix} = \begin{pmatrix} a \\ b \end{pmatrix}$  exists.  $\square$

If, moreover,  $V$  obeys (1.3) (a mild restriction), this and Theorem 1 imply that  $\sigma_{ac}(H) = [0, \infty)$ ,  $\sigma_{sing} \cap (0, \infty) = \emptyset$ .

#### NOTE ADDED IN PROOF

Subsequent to our work, S. Jitomirskaya and Y. Last [Phys. Rev. Lett. **76** (1996), 1765–1769 and papers in preparation] have shown how to get explicit  $m$ -function bounds in the general situation covered by Gilbert-Pearson theory.

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