# The Golinskii-Ibragimov Method and a Theorem of Damanik and Killip

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### 1 Introduction

Let  $\partial \mathbb{D}$  be the unit circle  $\{z \mid |z|=1\}$  in  $\mathbb{C}$  and  $\mathbb{D}$  the open disc  $\{z \mid |z|<1\}$ . Let  $\mu$  be a probability measure on  $\partial \mathbb{D}$  which is not supported on a finite number of points. Then using the Gram-Schmidt procedure, we can define monic orthogonal polynomials on the unit circle (OPUC)  $\Phi_n(z;d\mu)$  and normalized polynomials  $\phi_n(z;d\mu)$ . These obey the Szegö recursion formulae [9,6]:

$$\Phi_{n+1}(z) = z\Phi_n(z) - \bar{\alpha}_n\Phi_n^*(z), \tag{1.1}$$

where

$$\Phi_{n}^{*}(z) = z^{n} \overline{\Phi_{n}\left(\frac{1}{\bar{z}}\right)}. \tag{1.2}$$

The parameters  $\alpha_n$  are called the Verblunsky coefficients of  $d\mu$  (also called Schur, Szegö, Geronimus, or reflection parameters or coefficients). They lie in  $\mathbb{D}$ , and any  $\alpha \in \times_{n=0}^{\infty} \mathbb{D}$  is the Verblunsky coefficient of a unique measure [2, 9]. In this paper, we are mainly interested in the spectral problem of going from information on the Verblunsky coefficients to information on the measure.

Our starting point is a method from a lovely 1971 paper of Golinskii and Ibragimov [4], who used this method to prove the following theorem.

Received 13 March 2003. Communicated by Percy Deift. **Theorem 1.1** (Golinskii and Ibragimov [4]). If  $\sum_{n=0}^{\infty} n |\alpha_n|^2 < \infty$ , then  $d\mu_s = 0$ .

Here, we make a canonical decomposition:

$$d\mu = w(\theta) \frac{d\theta}{2\pi} + d\mu_s, \tag{1.3}$$

where  $d\mu_s$  is singular with respect to  $d\theta$ . We are interested in  $\alpha$ 's that obey the following weaker condition:

$$\sum_{n=0}^{N} n \left| \alpha_n \right|^2 \le A \log N + C \tag{1.4}$$

with A and C constants. Think of  $\alpha_n=\sqrt{A}/n$  as a prototypical example. First, we prove the following theorem.

**Theorem 1.2.** If (1.4) holds with 
$$A < 1/4$$
, then  $d\mu_s = 0$ .

The Golinskii-Ibragimov (GI) method most directly only gets A < 1/16, but by replacing their  $L^1$  method by an  $L^2$  method, we bring things up to A < 1/4.

Theorem 1.2 is almost optimal in that there are examples with A>1/4 but |A-1/4| arbitrarily small, where dµ has an eigenvalue. However, we can do better if we assume the  $\alpha$ 's are real.

Theorem 1.3. If all  $\alpha_n$  are real (equivalently,  $d\mu$  is invariant under complex conjugation) and (1.4) holds with A < 1/2, then  $d\mu_s$  can only consist of possible pure points at z = +1 or z = -1. If A = 1/4,  $d\mu_s = 0$ .

We use an idea motivated by Damanik and Killip [1] to prove Theorem 1.3. In this regard, we prove the following special looking result which, as we explain, is related to [1].

**Theorem 1.4.** Let all  $\alpha_n$  be real and obey the following conditions:

(i) the bound

$$\sum_{j=0}^{N} j \left| \alpha_{2j} \right|^2 \le A \log N + C, \tag{1.5}$$

- $(ii) \ |\alpha_{2j-1}| \leq |\alpha_{2j+1}| \ for \ j=1,2,\dots \ with \ either \ all \ \alpha_{2j-1} \geq 0 \ or \ all \ \alpha_{2j-1} \leq 0,$
- (iii)  $\sum_{i=1}^{\infty} |\alpha_{2i-1}|^2 < \infty$ .

If A<1/2,  $d\mu_s$  consists only of possible pure points at  $z=\pm 1$  or  $z=\pm i$ . Moreover, if A=1/4, the only possible pure points are at  $z=\pm 1$ , and if A=1/4 and  $\alpha_{2j-1}\leq 0$ , then  $d\mu_s=0$ .

This theorem is custom-made to provide part of a proof of the following recent striking result of Damanik and Killip [1].

**Theorem 1.5** (Damanik and Killip [1]). Let H be a half-line Schrödinger operator on  $\ell^2(\mathbb{Z}_+)$ ,

$$(\mathrm{Hu})_{\mathrm{n}} = \mathrm{u}_{\mathrm{n}+1} + \mathrm{u}_{\mathrm{n}-1} + \mathrm{v}_{\mathrm{n}} \mathrm{u}_{\mathrm{n}} \tag{1.6}$$

 $(u_0 \text{ interpreted as 0}). \text{ Suppose spec}(H) = [-2,2]. \text{ Then H has purely a.c. spectrum, that is,} \\ \sigma_{sc} = \sigma_{pp} = \varnothing. \\ \square$ 

Their proof has the following steps (they use  $\gamma_n$  for our  $\alpha_n$ ).

(i) Following Szegő [6], map H to an associated measure on  $\partial \mathbb{D}$  by using  $z=e^{i\theta}\mapsto E=2\cos\theta$  to pull back the spectral measure  $d\rho$  on [-2,2] for H to a measure  $\mu$  on  $\partial \mathbb{D}$ . Note (following Geronimus) that the Verblunsky coefficients for  $d\mu$  obey that

$$v_{n+1} = (1 - \alpha_{2n-1})\alpha_{2n} - (1 + \alpha_{2n-1})\alpha_{2n-2}, \tag{1.7}$$

$$1 = (1 - \alpha_{2n-1})(1 - \alpha_{2n}^2)(1 + \alpha_{2n+1})$$
(1.8)

(for general Jacobi matrices, the left-hand side of (1.8) is  $\alpha_{n+1}^2$ ; the initial conditions are  $\alpha_{-1}=-1$ ).

- (ii) Analyze (1.8) with  $\alpha_n$  real and  $|\alpha_n| < 1$  to conclude the following:
  - (a)  $\alpha_{2n-1} \leq \alpha_{2n+1} \leq 0$ , [1, Lemma 4.1],
  - (b)  $|\alpha_{2n+1}| \leq 1/(n+2)$  so  $\sum_{n=1}^{\infty} |\alpha_{2n-1}|^2 < \infty,$  [1, Lemma 4.2],
  - (c)  $\sum_{i=0}^{N} (j+1)\alpha_i^2 \leq (1/4)\log N + C, [1, Proposition 4.5].$
- (iii) Translate information on the  $\alpha_n$ 's to information on the  $\nu_n$ 's.
- (iv) Prove that -2, 2, 0 are not eigenvalues of H.
- (v) Prove that solutions of

$$u_{n+1} + u_{n-1} + v_n u_n = Eu \tag{1.9}$$

with  $E \in (-2,0) \cup (0,2)$  have  $|u_n| \le c n^\eta$  for any  $\eta > (1/2\sqrt{2})$ .

- (vi) Prove that the set of E in (-2,2), for which (1.9) has unbounded solutions, has Hausdorff dimension 0.
- (vii) Use (v), (vi), Hausdorff dimension, and the Jitomirskaya-Last inequalities [7] to show that  $d\mu$  has no singular part on  $(-2,0) \cup (0,2)$ .

Given step (ii), one can use Theorem 1.4 to replace steps (iii), (iv), (v), (vi), and (vii). For Theorem 1.4 says that  $d\mu$  is purely a.c. and the pullback then implies that  $d\rho$  is.

Since Theorem 1.4 depends on ideas closely related to steps (iv) and (v), what we are really doing is using an appeal to the GI method to replace steps (vi) and (viii) and, in particular, the use of Hausdorff dimension and the Jitomirskaya-Last inequalities. These steps follow the ideas of Remling [8] so, in essence, where [1] extends [8], we extend [4]. It is pointless to argue which approach is "simpler" (since some of their techniques have appeared extensively in the Schrödinger operator literature), but we believe that it is useful to have the alternate approach.

In Section 2, we discuss the ideas of Golinskii-Ibragimov and, in particular, prove Theorem 1.2. In Section 3, we use Prüfer variables for OPUC to prove Theorem 1.3. In Section 4, we prove Theorem 1.4.

#### 2 The GI method

Golinskii and Ibragimov made the assumption that

$$\sum_{n=0}^{\infty} n \left| \alpha_n \right|^2 < \infty. \tag{2.1}$$

On  $\partial \mathbb{D}$ ,  $D^{-1}$  is defined a.e.  $d\theta$  as boundary values of D, the Szegö function (see [3, 9]), and so a.e.  $d\mu_{ac}$ . We extend it to  $L^2(\partial \mathbb{D}, d\mu) = L^2(\partial \mathbb{D}, d\mu_{ac}) \oplus L^2(\partial \mathbb{D}, d\mu_s)$  by setting to 0 on the singular subspace. Then Golinskii and Ibragimov [4] proved that

$$\|\phi_{n}^{*} - D^{-1}\|_{L^{2}(\partial \mathbb{D}, d\mu)} \le Cn^{-1/2} \tag{2.2}$$

and that

$$\left\| \left( \phi_n^* \right)^{-1} \right\|_{\infty} \le C_1 \exp \left( C_2 \sqrt{\log n} \right), \tag{2.3}$$

where  $\|\cdot\|_{\infty}$  means sup over  $\partial\mathbb{D}$  or  $\bar{\mathbb{D}}$  (equal since  $(\phi_n^*)^{-1}$  is analytic in a neighborhood of  $\bar{\mathbb{D}}$ ). In (2.2), the Szegö function exists since (2.1) implies  $\sum_{n=0}^{\infty} |\alpha_n|^2 < \infty$ . We actually prove results like this below.

They then wrote

$$\begin{split} \left\| \left| \boldsymbol{\phi}_{n}^{*} \right|^{-2} - |D|^{2} \right\|_{L^{1}(\partial \mathbb{D}, d\theta/2\pi)} \\ &\leq \left\| \left( \boldsymbol{\phi}_{n}^{*} \right)^{-1} \right\|_{\infty}^{2} \left\| \left( \left| \boldsymbol{\phi}_{n} \right|^{2} - |D|^{-2} \right) D^{2} \right\|_{L^{1}(\partial \mathbb{D}, d\theta/2\pi)} \end{split} \tag{2.4}$$

$$\leq \left\| \left(\phi_{\mathfrak{n}}^{*}\right)^{-1} \right\|_{\infty}^{2} \left\| \left( \left|\phi_{\mathfrak{n}}\right|^{2} - |D|^{-2} \right) \right\|_{L^{1}(\partial \mathbb{D}; d\mu)} \tag{2.5}$$

$$\leq \left\| \left(\phi_{\mathfrak{n}}^{*}\right)^{-1} \right\|_{\infty}^{2} \left\| \phi_{\mathfrak{n}} - D^{-1} \right\|_{L^{2}(\partial \mathbb{D}; d\mu)} \left\| \left| \phi_{\mathfrak{n}} \right| + |D|^{-1} \right\|_{L^{2}(\partial \mathbb{D}; d\mu)}. \tag{2.6}$$

In (2.5), we use  $D^2(d\theta/2\pi) = d\mu_{ac} \le d\mu$  and in (2.6) we use  $||\phi_n| - |D|^{-1}| \le |\phi_n - D^{-1}|$ . Thus, by (2.2) and (2.3),  $|\phi_n|^{-2}(d\theta/2\pi) \to |D|^2(d\theta/2\pi)$  in norm on measures. But it is known that  $|\phi_n|^{-2}(d\theta/2\pi) \to d\mu$  weakly (see, e.g., [2]). Thus  $d\mu = |D|^2(d\theta/2\pi)$  so  $d\mu_s = 0$ .

What is especially interesting about this approach is that it uses the divergent estimate (2.3). Clearly, we can have much more rapid growth of  $\|(\phi_n^*)^{-1}\|_\infty$  than in (2.3) and still have convergence. Basically, it suffices to have  $\|(\phi_n^*)^{-1}\|_\infty \leq C n^\beta$ , with  $2\beta < 1/2$ . If one looks at the proof of (2.3) in [4], that translates to a bound like (1.4) with  $\sqrt{A} < 1/4$ , that is, A < 1/16. Our first observation is that instead of estimating L¹-norms as GI do, it pays to estimate L²-norms. Then  $\|(\phi_n^*)^{-1}\|_\infty$  will occur as a first power, not second. Here are the key facts.

**Theorem 2.1.** Let μ obey the Szegö condition. Then

- (a)  $d\mu_s = 0$  if and only if  $\|(\phi_n^*)^{-1} D\|_{L^2(\partial \mathbb{D}, d\theta/2\pi)} \to 0$ ;
- $(b) \ \ \text{if I is an open interval and} \ \|\chi_I[(\phi_n^*)^{-1}-D]\|_{L^2(\partial\mathbb{D},d\theta/2\pi)} \to 0, \text{then } \mu_s(I)=0.$

Proof. (a) As is well known (see [2]),  $\|(\phi_n^*)^{-1}\|_{L^2(\partial \mathbb{D}, d\theta/2\pi)} = 1$ . Moreover,  $(\phi_n^*)^{-1}(z) \to D(z)$  uniformly on compact subsets of  $\mathbb{D}$ , so  $(\phi_n^*)^{-1} \to D$  weakly in  $L^2$ . Thus  $(\phi_n^*)^{-1} \to D$  in norm if and only if  $1 = \|(\phi_n)^{-1}\|_{L^2(\partial \mathbb{D}, d\theta/2\pi)}^2 = \|D\|_{L^2(\partial \mathbb{D}, d\theta/2\pi)}^2$ . But since  $\|D\|^2 = \int |D(\theta)|^2 (d\theta/2\pi) = \mu_{ac}(\partial \mathbb{D})$ , norm convergence is equivalent to  $\mu_s(\partial \mathbb{D}) = 0$ .

(b) Let f be continuous and have support in I. Then

$$f\big|\phi_n^*\big|^{-2} = \Big(f\big(\phi_n^*\big)^{-1}\chi_1\Big)\Big(\overline{\big(\phi_n^*\big)^{-1}}\chi_I\Big) \longrightarrow f|D|^2 \tag{2.7}$$

in  $L^1$  and thus, by the weak convergence of  $|\phi_n^*|^{-2}(d\theta/2\pi)$  to  $d\mu,$  we have

$$\int f \, d\mu = \int f |D|^2 \left( \frac{d\theta}{2\pi} \right), \tag{2.8}$$

that is, 
$$d\mu_s(I) = 0$$
.

Theorem 2.2. If

$$\left\|\chi_{I}\left(\phi_{n}^{*}\right)^{-1}\right\|_{\infty}\left\|\left(\phi_{n}^{*}-D^{-1}\right)\right\|_{L^{2}\left(\eth\mathbb{D},d\mu\right)}\longrightarrow0\tag{2.9}$$

for some open interval I (including  $I = \partial \mathbb{D}$ ), then  $\mu_s(I) = 0$ .

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Proof. We have that

$$\begin{split} \left\| \chi_{\rm I} \Big[ \big( \phi_n^* \big)^{-1} - D \Big] \right\|_{L^2(\partial \mathbb{D}, d\theta/2\pi)} &= \left\| \chi_{\rm I} \big( \phi_n^* \big)^{-1} \big( \phi_n^* - D^{-1} \big) D \right\|_{L^2(\partial \mathbb{D}, d\theta/2\pi)} \\ &\leq \left\| \chi_{\rm I} \big( \phi_n^* \big)^{-1} \right\|_{\infty} \left\| \big( \phi_n^* - D^{-1} \big) \right\|_{L^2(\partial \mathbb{D}, d\mu_{ac})} \\ &\leq LHS \ of \ (2.9). \end{split} \tag{2.10}$$

Thus the result follows from Theorem 2.1.

Define  $\rho_{\mathfrak{n}} = (1-|\alpha_{\mathfrak{n}}|^2)^{1/2}.$  A straightforward calculation shows that for  $\mathfrak{n} < \mathfrak{m},$ 

$$\left\langle \phi_n^*, \phi_m^* \right\rangle = \prod_{\ell=n}^{m-1} \rho_\ell \tag{2.11}$$

and, of course,  $\langle \phi_n^*, \phi_n^* \rangle = 1$ . Since [3]  $\phi_n^* \to D^{-1}$  in  $L^2(\partial \mathbb{D}, d\mu)$ , we have that

$$\|\varphi_{n}^{*} - D^{-1}\|_{L^{2}(\partial \mathbb{D}, d\mu)}^{2} = 2\left(1 - \prod_{\ell=n}^{\infty} \rho_{\ell}\right). \tag{2.12}$$

 $\text{Proposition 2.3. (a) We have that } \|\phi_{\mathfrak{n}}^* - D^{-1}\|_{L^2(\mathfrak{dD}, d\mu)}^2 \leq 2 \sum_{\ell=n}^\infty |\alpha_\ell|^2.$ 

- (b) If (1.4) holds, then  $\sum_{n=0}^{\infty} |\alpha_n|^2 < \infty.$
- (c) If (1.4) holds for any A, then

$$\|\phi_{n}^{*} - D^{-1}\|_{L^{2}(\partial \mathbb{D}, d\mu)} \le \frac{C \log n}{n^{1/2}}$$

for all large n.

Proof. (a) Since  $\rho_\ell^2 \leq \rho_\ell, 1 - \prod_{\ell=n}^\infty \rho_\ell \leq 1 - \prod_{\ell=n}^\infty \rho_\ell^2$ . Moreover,  $1 - \prod_{j=1}^\ell x_j \leq \sum_{j=1}^\ell (1 - x_j)$  by a simple induction. Thus (2.3) implies (a), (b), (c)

$$\sum_{n=2^{\ell}}^{2^{\ell+1}-1} \left|\alpha_{n}\right|^{2} \leq 2^{-\ell} \sum_{n=0}^{2^{\ell+1}} n \left|\alpha_{n}\right|^{2} \leq 2^{-\ell} \big[ A \log \big(2^{\ell+1}\big) + C \big]. \tag{2.14}$$

Since

$$\sum_{\ell=k}^{\infty} \ell 2^{-\ell} \le Ck 2^{-k},\tag{2.15}$$

we see that (1.4) implies  $\sum_{n=0}^{\infty} |\alpha_n|^2 < \infty$  and

$$\sum_{\ell=n}^{\infty} \left| \alpha_{\ell} \right|^2 \le C n^{-1} \log n. \tag{2.16}$$

Given (a), we get (c).

**Theorem 2.4.** If the  $\alpha$ 's obey (1.4) and for some open interval  $I \subset \partial \mathbb{D}$ , and B < 1/2, then

$$\sup_{z \in I} \left[ \left| \varphi_{\mathfrak{n}}(z) \right|^{-1} \right] \le C(\mathfrak{n} + 1)^{B} \tag{2.17}$$

for a constant C, then  $\mu_s(I) = 0$ .

Proof. For  $z \in \partial \mathbb{D}$ ,  $|\phi_n^*(z)| = |\phi_n(z)|$  so (2.17) states that  $\|\chi_I(\phi_n^*)^{-1}\|_{\infty} \leq C(n+1)^B$ . This and (2.13) imply that (2.9) holds, so  $\mu_s(I) = 0$ .

The second kind polynomials,  $\psi_n$  and  $\Psi_n$ , are defined by reversing the signs of all the Verblunsky coefficients. It is known (see, e.g., [5]) that for  $z \in \partial \mathbb{D}$ ,  $\text{Re}(\phi_n(z)\bar{\psi}_n(z)) = 1$ , so  $|\phi_n(z)|^{-1} \leq |\psi_n(z)|$ .

Since  $\psi_n = (\prod_{j=0}^{n-1} \rho_j^{-1}) \Psi_n$  and  $\sum \alpha_j^2 < \infty$  implies  $\prod_{j=0}^{\infty} \rho_j^{-1} < \infty$ , we see that  $|\psi_n(z)| \le C|\Psi_n|$ . Thus, Theorem 2.4 can be rewritten as the following corollary.

**Corollary 2.5.** If the  $\alpha$ 's obey (1.4) and for some open interval  $I \subset \partial \mathbb{D}$ , we have

$$\sup_{z \in I} \left| \Psi_n(z) \right| \le C(n+1)^B \tag{2.18}$$

for a constant C, and B < 1/2, then  $\mu_s(I) = 0$ .

Proof of Theorem 1.2. Equation (1.2) implies  $|\Psi_n^*| = |\Psi_n|$  on  $\partial \mathbb{D}$  and so (1.1) implies  $|\Psi_{n+1}| \le (1 + |\alpha_n|) |\Psi_n|$ . Thus

$$\sup_{z\in\partial\mathbb{D}}\left|\Psi_{n}(z)\right|\leq\prod_{j=0}^{n-1}1+\left|\alpha_{j}\right|\leq\exp\left(\sum_{j=0}^{n-1}\left|\alpha_{j}\right|\right).\tag{2.19}$$

But, by the Schwartz inequality, (C is a "variable constant" and  $\varepsilon$  is arbitrarily small)

$$\begin{split} \sum_{j=0}^{n} \left| \alpha_{j} \right| &\leq \left( \sum_{j=0}^{n} (j+1) \left| \alpha_{j} \right|^{2} \right)^{1/2} \left( \sum_{j=0}^{n} (j+1)^{-1} \right)^{1/2} \\ &\leq \left[ A \log(n) + C \right]^{1/2} \left[ \log(n) + C \right]^{1/2} \\ &\leq \sqrt{A + \varepsilon} \log(n) + C. \end{split} \tag{2.20}$$

Thus, by (2.19),

$$\sup_{z \in \mathbb{D}} \left| \Psi_{n}(z) \right| \le C n^{\sqrt{A+\varepsilon}}. \tag{2.21}$$

If  $\sqrt{A + \varepsilon} < 1/2$ , that is, A < 1/4, Corollary 2.5 is applicable.

We emphasize that, in essence, the calculation in Proposition 2.3(a), (2.19), and the basic strategy are all from [4]; the only real advance in this section is the use of  $L^2$ -norms allowing A < 1/4, where the method of [4] gets A < 1/16.

Finally, we note that A=1/4 is a critical value for the appearance of bound states, on account of the following theorem.

**Theorem 2.6.** If all  $\alpha_n$  are real, then  $\phi_n(1) > 0$  and

$$\varphi_{n}(1) = \prod_{j=0}^{n-1} (1 - \alpha_{j}).$$
(2.22)

In addition,  $(-1)^n \phi_n(-1) > 0$  and

$$(-1)^{n} \varphi_{n}(-1) = \prod_{j=0}^{n-1} \left(1 - (-1)^{j+1} \alpha_{j}\right). \tag{2.23}$$

Remark 2.7. In particular, if  $\alpha_j=B(j+1)^{-1}$ , then  $\phi_n(1)\sim n^{-B}$  and  $\sum |\phi_n(1)|^2<\infty$  if B>1/2. For that  $\alpha$ , (1.4) holds for  $A=B^2$ , that is, there are examples with  $d\mu_s\neq 0$  and (1.4) holding for any A>1/4. At A=1/4, our proof of Theorem 1.2 shows that  $|\alpha_n(z)|\geq Cn^{-1/2}$ , so at least there are no eigenvalues.

Proof. Note that  $\varphi_n(1)$  is real by induction and then (1.1) states that

$$\phi_{n+1}(1) = (1 - \alpha_n)\phi_n(1), \tag{2.24}$$

proving (2.22). Similarly,  $\varphi_n(-1)$  is real, (1.1) for z = -1 states that

$$(-1)^{n+1}\phi_{n+1}(-1) = (-1)^n\phi_n(-1) - \alpha_n(-1)^{n+1}(-1)^n\phi_n(-1) \tag{2.25}$$

since 
$$\varphi_n^*(-1) = (-1)^n \varphi_n(-1)$$
, which yields (2.23).

# 3 Prüfer variables and the proof of Theorem 1.3

Write  $\Psi_n(z)=R_n(z)e^{i[\theta_n(z)+n\eta(z)]}$  where  $e^{i\eta(z)}=z$  and  $\theta_n$  is defined initially only modulo  $2\pi$ . If  $\alpha_n$  is real, (1.1) with  $\alpha_n\to -\alpha_n$  becomes  $R_{n+1}^2=R_n^2|e^{i\eta}e^{i[\theta_n+n\eta]}+\alpha_ne^{-i\theta_n}|^2$ . Thus,

$$\frac{R_{n+1}^2}{R_n^2} = 1 + \alpha_n^2 + 2\alpha_n \cos((n+1)\eta + 2\theta_n). \tag{3.1}$$

Secondly,

$$e^{i(\theta_{n+1}(z)-\theta_n(z))} = \frac{1 + \alpha_n e^{-i(2\theta_n + (n+1)\theta_n)}}{\left[1 + \alpha_n^2 + 2\alpha_n \cos\left((n+1)\eta + 2\theta_n\right)\right]^{1/2}}.$$
 (3.2)

These are the Prüfer variable equations for  $\Psi$ , (3.2) implies that  $\cos(\theta_{n+1}-\theta_n)>0$ , so we can always pick  $\theta_{n+1}$  so  $|\theta_{n+1}-\theta_n|<\pi/2$ , settling the  $2\pi$  ambiguity.

**Proposition 3.1.** Let  $\alpha_n$  be real and  $\sum \alpha_n^2 < \infty$ . Define  $Q = \sup_n |\alpha_n| < 1$ . Then

(a) for a z independent constant  $C \in (1, \infty)$ ,

$$C^{-1} \leq R_n \, exp \left( - \sum_{j=0}^{n-1} \alpha_j \cos \left( (j+1) \eta + 2\theta_j(z) \right) \right) \leq C, \tag{3.3}$$

(b) for all n,

$$\left|\theta_{n+1} - \theta_n\right| \le \frac{\pi}{2} \frac{1}{1 - Q} \left|\alpha_n\right|. \tag{3.4}$$

Remark 3.2. The constant C depends only on Q and  $\sum_{j=0}^{\infty} \alpha_j^2$ .

Proof. (a) Define  $b_n(z)$  so the right-hand side of (3.1) is  $1 + b_n$ . Then

$$(1-Q)^2 \le (1-|\alpha_n|)^2 \le 1 + b_n \le (1+\alpha_n)^2 \le (1+Q)^2. \tag{3.5}$$

It follows that for Q dependent constant K, we have  $e^{-Kb_\pi^2} \le (1+b_\pi)e^{-b_\pi} \le e^{Kb_\pi^2}$ . Since  $b_n \le 3|\alpha_n|$  and  $b_n$  has an  $\alpha_n^2$  in it, we have

$$\begin{split} \exp\left(-\left(9K+1\right)\alpha_{n}^{2}\right) &\leq \left(1+b_{n}\right)\exp\left(-2\alpha_{n}\cos\left((n+1)\eta+2\theta_{n}\right)\right) \\ &\leq \exp\left(\left(9K+1\right)\alpha_{n}^{2}\right). \end{split} \tag{3.6}$$

Thus, by (3.1) and  $R_0 = 1$ , we have (3.3) with

$$C = \exp\left((9K + 1)\sum_{j=0}^{\infty} \alpha_j^2\right). \tag{3.7}$$

(b) Taking imaginary parts of both sides of (3.2) and using the lower bound 1-Q on the denominator, we get

$$\left|\sin\left(\theta_{n+1} - \theta_n\right)\right| \le \left|\alpha_n\right| (1 - Q)^{-1}. \tag{3.8}$$

Equation (3.2) also implies that  $\cos(\theta_{n+1}-\theta_n)>(1-Q)/(1+Q)>0$ , so  $|\theta_{n+1}-\theta_n|<\pi/2$ . Since  $|x_n|<\pi/2$  implies  $|x|\leq (\pi/2)|\sin x|$ , (3.4) follows from (3.8).

The point of (3.3) is to control  $|R_n|$ , we need to control  $\sum_{j=0}^{n-1} \alpha_j \cos((j+1)\eta + 2\theta_j)$ . In using the Schwartz inequality, we decouple the  $\alpha_j$ 's and the cosines so the key will be the following lemma (essentially in [1]).

**Lemma 3.3.** (i) If k is not a multiple of  $2\pi$ ,

$$\left|\sum_{j=1}^{n} \frac{\cos\left(kj + \theta_{j}\right)}{j}\right| \leq \frac{1}{\left|\sin\left(\frac{k}{2}\right)\right|} \left[1 + \sum_{j=1}^{n-1} \frac{\left|\theta_{j+1} - \theta_{j}\right|}{j+1}\right]. \tag{3.9}$$

(ii) If k is not a multiple of  $\pi$ ,

$$\sum_{j=1}^{N} \frac{\cos^2 \left( kj + \theta_j \right)}{j} \le \frac{1}{2} (\log N + 1 + C), \tag{3.10}$$

where

$$C = \frac{1}{|\sin(k)|} \left[ 1 + 2 \sum_{j=1}^{\infty} \frac{|\theta_{j+1} - \theta_j|}{j+1} \right].$$
 (3.11)

Proof. (i) Since  $cos(kj+\theta_j)=Re(exp(ikj)\,exp(i\theta_j))$ , it suffices to prove (3.9) with cosines replaced by complex exponentials. Define  $b_n=\sum_{j=1}^n e^{ikj}$  so, by summing the geometric series,

$$\left|b_{n}\right| \leq \frac{1}{\left|\sin\left(\frac{k}{2}\right)\right|}.\tag{3.12}$$

If  $a_j=e^{i\theta_j}/j$ , then the sum we want to control is  $\sum_{j=1}^n(b_j-b_{j-1})a_j$  with  $b_0=0$ . But

$$\sum_{j=1}^{n} (b_{j} - b_{j-1}) a_{j} = -\sum_{j=1}^{n} b_{j} (a_{j+1} - a_{j}), \tag{3.13}$$

where  $a_{n+1} = 0$ . Thus

$$\left| \sum_{j=1}^{n} \frac{\cos\left(kj + \theta_{j}\right)}{j} \right| \leq \frac{1}{\sin\left(\frac{k}{2}\right)} \left[ \sum_{j=1}^{n-1} \left| \alpha_{j+1} - \alpha_{j} \right| + \left| \alpha_{n} \right| \right]. \tag{3.14}$$

Clearly,

$$\begin{split} \left| \alpha_{j+1} - \alpha_{j} \right| &\leq \left| \left( e^{i\theta_{j+1}} - e^{i\theta_{j}} \right) (j+1)^{-1} \right| + \left| e^{i\theta_{j}} \left( j^{-1} - (j+1)^{-1} \right) \right| \\ &\leq \left| \theta_{j+1} - \theta_{j} \right| (j+1)^{-1} + j^{-1} - (j+1)^{-1}. \end{split} \tag{3.15}$$

Since  $\sum_{j=1}^{n-1} j^{-1} - (j+1)^{-1} + n^{-1} = 1$ , (3.14) and (3.15) yield (3.9). (ii) Since  $\cos^2(x) = (1/2)(1 + \cos(2x))$  and

$$\sum_{i=1}^{n} \frac{1}{i} \le 1 + \int_{1}^{n} \frac{\mathrm{d}x}{x} = 1 + \log(n),\tag{3.16}$$

Proof of Theorem 1.3. Write

$$\begin{split} &\left| \sum_{j=0}^{n-1} \alpha_{j} \cos \left( (j+1)\eta + 2\theta_{j} \right) \right| \\ &\leq \left( \sum_{j=0}^{n-1} (j+1) \left| \alpha_{j} \right|^{2} \right)^{1/2} \left( \sum_{j=0}^{n-1} \frac{\cos^{2} \left( (j+1)\eta + 2\theta_{j} \right)}{(j+1)} \right)^{1/2}. \end{split} \tag{3.17}$$

By hypothesis, the first sum on the right-hand side of (3.17) is bounded by A log N + C. By the lemma, if  $\eta$  is not a multiple of  $\pi$  (i.e.,  $z \neq \pm 1$ ), the second sum is bounded by

$$\frac{1}{2} \Big( \log n + \widetilde{C} \big| \sin(\eta) \big|^{-1} \Big), \tag{3.18}$$

where C can be chosen independently of z, since

$$\begin{split} \sum_{j=1}^{\infty} \frac{\left| \theta_{j+1} - \theta_{j} \right|}{j+1} &\leq C \sum_{j=0}^{\infty} \frac{\left| \alpha_{j} \right|}{j+2} \quad (by \ (3.8)) \\ &\leq C \left( \sum_{j=0}^{\infty} \left| \alpha_{j} \right|^{2} \right)^{1/2} \left( \sum_{j=0}^{\infty} \frac{1}{(j+2)^{2}} \right)^{1/2} \end{split} \tag{3.19}$$

is finite. Thus, for  $\eta \in [\theta_0, \pi - \theta_0],$ 

$$\begin{aligned} \left| R_{n} \right| &\leq C \exp \left| \left[ A (\log n) + C \right] \right|^{1/2} \left[ \frac{1}{2} \log n + \widetilde{C} \left| \sin \theta_{0} \right|^{-1} \right]^{1/2} \\ &\leq C_{\varepsilon} n^{\sqrt{(1/2)A + \varepsilon}}, \end{aligned} \tag{3.20}$$

where  $C_\epsilon$  depends on  $\epsilon$  and  $\theta_0$ . So long as  $\sqrt{(1/2)A} < 1/2$ , we can apply Corollary 2.5, that is, A < 1/2. We conclude  $d\mu_s(I) = 0$  for  $I = \pm(\theta_0, \pi - \theta_0)$ , that is,  $d\mu_s$  is supported on  $\{\pm 1\}$ . As already noted, at  $\pm 1$ ,  $|\phi_n(\pm 1)| \geq n^{-\sqrt{A}}$ , so if A = 1/4,  $\phi_n(\pm 1)$  are not in  $L^2$  and  $d\mu_s = 0$ .

# 4 Sequences of bounded variation and the proof of Theorem 1.4

To obtain Theorem 1.4, we need one more summation-by-parts argument that will supplement Lemma 3.3.

**Lemma 4.1.** If k is not a multiple of  $2\pi$ ,

$$\begin{split} &\left| \sum_{j=1}^{n} c_{j} \cos \left( kj + \theta_{j} \right) \right| \\ &\leq \frac{1}{\sin \left( \frac{k}{2} \right)} \left[ \sum_{j=1}^{\infty} \left| c_{j+1} - c_{j} \right| + \sup_{n} \left| c_{n} \right| + \sum_{j=1}^{\infty} \left| c_{j+1} \right| \left| \theta_{j+1} - \theta_{j} \right| \right]. \end{split} \tag{4.1}$$

Proof. As in the proof of Lemma 3.3, let  $b_n = \sum_{j=1}^n e^{ikj}$ , so (3.12) holds and  $a_k = c_j e^{i\theta_j}$ . Then summing by parts as in the earlier lemma,

$$\left| \sum_{j=1}^{n} c_{j} \cos \left( kj + \theta_{j} \right) \right| \leq \frac{1}{\sin \left( \frac{k}{2} \right)} \left[ \left| c_{n} \right| + \sum_{j=1}^{n-1} \left| a_{j+1} - a_{j} \right| \right]. \tag{4.2}$$

But

$$\left| a_{j+1} - a_{j} \right| \le \left| a_{j+1} \right| \left| e^{i\theta_{j+1}} - e^{i\theta_{j}} \right| + \left| c_{j+1} - c_{j} \right|, \tag{4.3}$$

Proof of Theorem 1.4. By (3.3) and the hypothesis that  $\sum_{j=0}^{\infty}|\alpha_{j}|^{2}<\infty,$ 

$$\left|\Psi_{n}(z)\right| \leq C \exp\left(\sum_{j=0}^{n-1} \alpha_{j} \cos\left((j+1)\eta + 2\theta_{j}\right)\right). \tag{4.4}$$

Write

$$\sum_{i=0}^{n-1} \alpha_{i} \cos ((i+1)\eta + 2\theta_{i}) = O_{n} + E_{n}, \tag{4.5}$$

where  $O_n$  is the sum over odd values of j and  $E_n$  over even values.

By Lemma 4.1,

$$\begin{aligned} \left| O_{n} \right| &\leq \frac{1}{\left| \sin \left( \frac{\eta}{2} \right) \right|} \left[ \sup_{n} \left| \alpha_{2n-1} \right| + \sum_{n=1}^{\infty} \left| \alpha_{2n+1} - \alpha_{2n-1} \right| \right. \\ &+ \left. \sum_{n=1}^{\infty} \left| \alpha_{2n+1} \right| \left| \theta_{2n+1} - \theta_{2n-1} \right| \right]. \end{aligned} \tag{4.6}$$

Since  $\alpha_{2n-1}$  is monotone,  $\sum_{n=1}^{\infty} |\alpha_{2n+1} - \alpha_{2n-1}| = |\alpha_1|$ . By (3.4),  $|\theta_{2n+1} - \theta_{2n-1}| \le C(|\alpha_{2n-1}| + |\alpha_{2n}|)$  so, since  $\sum_{n=0}^{\infty} |\alpha_n|^2 < \infty$ , (4.6) implies

$$\left|O_{n}\right| \le C \left|\sin\left(\frac{\eta}{2}\right)\right|^{-1}.\tag{4.7}$$

For  $E_n$ , we use the same argument as in the proof of Theorem 1.3, taking into account that the change in frequency from 2n to 2n+2 is  $2\eta+(\theta_{2n+2}-\theta_{2n})$ . Thus  $|\sin(\eta)|^{-1}$  becomes  $|\sin(2\eta)|^{-1}$ , and we find that

$$\left| \mathsf{E}_{\mathrm{n}} \right| \le (\mathsf{A} \log \mathsf{N} + \mathsf{C})^{1/2} \left( \frac{1}{2} \log \mathsf{N} + \mathsf{C} \big| \sin(2\eta) \big|^{-1} \right)^{1/2}.$$
 (4.8)

Inequalities (4.7) and (4.8) imply that for any  $\theta_0 > 0$  for all  $z \in (\theta_0, \pi/2 - \theta_0) \cup (\pi/2 + \theta_0, \pi - \theta_0)$ ,

$$\left|\Psi_{n}\right| \leq C_{\theta_{0},\varepsilon} n^{\varepsilon + \sqrt{A/2}} \tag{4.9}$$

which, by Corollary 2.5, implies  $\mu_s$  is restricted to  $\pm 1$  and  $\pm i$ .

To obtain the result on eigenvalues when A=1/4, note first that since  $|E_n| \leq \sum_{m=1}^{[n/2]} |\alpha_{2m}|$ , uniformly in z,

$$e^{|E_n|} \le Cn^{1/2},$$
 (4.10)

$$\begin{split} \text{if } A = 1/4. \text{ At } z = \pm i, (4.7) \text{ implies } e^{|O_{\pi}|} \text{ is bounded, so } |\phi_{\pi}(\pm i)| \geq C \pi^{-1/2} \text{ is not in } L^2. \\ \text{At } \pm 1, \text{ we use Theorem 2.6. Since } \alpha_{2j-1} \leq 0, (\textbf{2.22}) \text{ implies that} \end{split}$$

$$\left|\phi_{\mathfrak{n}}(\pm 1)\right| \geq \prod_{\substack{\mathfrak{j}=0\\\mathfrak{j} \text{ even}}}^{\mathfrak{n}} \left(1 - \left|\alpha_{\mathfrak{j}}\right|\right) \geq C \mathfrak{n}^{-1/2} \tag{4.11}$$

since  $1 - \alpha_{2j-1} \ge 1$ . So  $\pm 1$  are not eigenvalues if  $\alpha_{2j-1} \le 0$ .

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### References

- [1] D. Damanik and R. Killip, Half-line Schrödinger operators with no bound states, preprint, 2003.
- [2] T. Erdélyi, P. Nevai, J. Zhang, and J. S. Geronimo, A simple proof of "Favard's theorem" on the unit circle, Atti Sem. Mat. Fis. Univ. Modena 39 (1991), no. 2, 551–556.
- [3] G. Freud, Orthogonal Polynomials, Pergamon Press, Oxford, 1971.
- [4] B. L. Golinskii and I. A. Ibragimov, On Szegö's limit theorem, Math. USSR-Izv. 5 (1971), 421–444.
- [5] L. Golinskii and P. Nevai, Szegö difference equations, transfer matrices and orthogonal polynomials on the unit circle, Comm. Math. Phys. 223 (2001), no. 2, 223–259.
- [6] G. Szegö, *Orthogonal Polynomials*, 3rd ed., American Mathematical Society Colloquium Publications, vol. 23, American Mathematical Society, Rhode Island, 1967.
- [7] S. Jitomirskaya and Y. Last, *Power-law subordinacy and singular spectra*. *I. Half-line operators*, Acta Math. **183** (1999), no. 2, 171–189.
- [8] C. Remling, The absolutely continuous spectrum of one-dimensional Schrödinger operators with decaying potentials, Comm. Math. Phys. 193 (1998), no. 1, 151–170.
- [9] B. Simon, *Orthogonal Polynomials on the Unit Circle*, to appear in AMS Colloquium Publication Series.

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