



Gap Labelling for Periodic Jacobi Matrices on Trees

Barry Simon

IBM Professor of Mathematics and Theoretical Physics, Emeritus
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Pasadena, CA, U.S.A.

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Joint Work with Jess Banks (Berkeley), Jonathan Breuer (HUJI), Jorge
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After discussing periodic Jacobi matrices on trees and this theorem, I'll discuss the more general context of gap labelling and the historical model of Floquet theory for Hill's equation. Then after describing the tools we'll need, I'll focus on a miraculous formula that will prove Sunada's theorem.



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After discussing periodic Jacobi matrices on trees and this theorem, I'll discuss the more general context of gap labelling and the historical model of Floquet theory for Hill's equation. Then after describing the tools we'll need, I'll focus on a miraculous formula that will prove Sunada's theorem. If there is time I'll discuss another result for which our new formula provides a new proof.



Graphs

Let Γ be a graph with vertex set, $\mathbb{V}(\Gamma)$, an edge set, $\mathbb{E}(\Gamma)$.

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Let Γ be a graph with vertex set, $\mathbb{V}(\Gamma)$, an edge set, $\mathbb{E}(\Gamma)$. We will suppose some familiarity with notions of graph theory but will remind about some terminology. In particular, we recall that the *degree* of a vertex, $v \in \mathbb{V}(\Gamma)$ is the number of edges, $e \in \mathbb{E}(\Gamma)$, with v as an endpoint and that a *leaf* is a vertex of degree 1.

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We will assign an orientation for each edge, e , using \check{e} for the oppositely directed edge. $\sigma(e)$ is the initial vertex and $\tau(e)$ the final of the directed edge e , so for example, $\sigma(\check{e}) = \tau(e)$. We let $\tilde{\mathbb{E}}$ denote the set of all edges with arbitrary assigned orientation so that $\#\tilde{\mathbb{E}} = 2\#\mathbb{E}$.



Jacobi Matrices on Graphs

A *Jacobi matrix* on Γ is defined by *Jacobi parameters*, i.e. a potential, $b(v) \in \mathbb{R}$, to each vertex and coupling, $a(e) = a(\check{e}) > 0$, to each edge.

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The Jacobi matrix is indexed by pairs of vertices and defines an operator on $\ell^2(\mathbb{V}(\Gamma))$ by taking

$$H_{vw} = \begin{cases} b(v), & \text{if } v = w \\ a(e), & \text{if } (vw) = e \text{ an edge in } \tilde{\mathbb{E}}(\Gamma) \\ 0, & \text{otherwise} \end{cases}$$

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If the graph has n -vertices, $\{1, \dots, n\}$ with edges between j and $j + 1$, this is a classical tridiagonal Jacobi matrix.

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If the graph has n -vertices, $\{1, \dots, n\}$ with edges between j and $j + 1$, this is a classical tridiagonal Jacobi matrix. If the graph has vertex set \mathbb{Z} with neighboring edges, we get a classical (doubly) infinite Jacobi matrix.

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Periodic Jacobi Matrices on Trees

Now let Γ be a finite, leafless graph. Such a graph always has loops, i.e. is not simply connected.

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Periodic Jacobi Matrices on Trees

Now let Γ be a finite, leafless graph. Such a graph always has loops, i.e. is not simply connected. If Γ has ℓ independent loops (equivalently, one can drop at most ℓ edges without disconnecting the graph), the fundamental group is \mathbb{F}_ℓ , the nonabelian free group on ℓ generators (which despite the name is abelian if (and only if) $\ell = 1$).

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Let \mathcal{T} be the universal cover of Γ . It is a tree.

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Let \mathcal{T} be the universal cover of Γ . It is a tree. There is a cover map $\pi : \mathcal{T} \rightarrow \Gamma$ and a family of deck transformations isomorphic to \mathbb{F}_ℓ which acts transitively on each $\pi^{-1}(v)$ for each $v \in \mathbb{V}(\Gamma)$. Given a Jacobi matrix, J_Γ , on Γ , with Jacobi parameters, b and a , there is a unique lift to Jacobi parameters on \mathcal{T} given by $b(\tilde{v}) = b(\pi(\tilde{v}))$, $a(\tilde{e}) = a(\pi(\tilde{e}))$.



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Periodic Jacobi Matrices on Trees

Two simple and canonical examples are where first Γ is a single cycle with p vertices and second where Γ has a two vertices with d edges connecting them.

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Periodic Jacobi Matrices on Trees

Two simple and canonical examples are where first Γ is a single cycle with p vertices and second where Γ has two vertices with d edges connecting them. In the first case $H_{\mathcal{T}}$ is a conventional periodic Jacobi matrix of period p (which is where our notion of period comes from) - a subject on which there is truly enormous history and literature which we will discuss in part below.

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Periodic Jacobi Matrices on Trees

In the second, \mathcal{T} is a homogenous tree of degree d , where each vertex has degree d . This is interesting even in case all b are 0 and a 's are equal.

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In the second, \mathcal{T} is a homogenous tree of degree d , where each vertex has degree d . This is interesting even in case all b are 0 and a 's are equal. As constructed, it has period 2.

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In the second, \mathcal{T} is a homogenous tree of degree d , where each vertex has degree d . This is interesting even in case all b are 0 and a 's are equal. As constructed, it has period 2. Related is the case where Γ has one vertex and ℓ self loops. In that case, the tree is homogeneous of degree 2ℓ ; $H_{\mathcal{T}}$ is a special case of the last class where the 2ℓ values of a occur in ℓ pairs (and b 's are all the same). Thus some of that second class are “secretly” period 1.



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The model of the homogeneous tree and these models more generally are connected to modular forms and so this subject is of interest to mathematical physicists, spectral theorists and number theorists.



The DOS and Gap Labelling

Deck transformations induce unitary maps on $\ell^2(\mathbb{V}(\mathcal{T}))$ which commute with $H_{\mathcal{T}}$. In particular, for every $v \in V(\Gamma)$, the spectral measure, $d\mu_{\tilde{v}}$, for $H_{\mathcal{T}}$ and \tilde{v} are the same for all $\tilde{v} \in \mathbb{V}(\mathcal{T})$ with $\pi(\tilde{v}) = v$.

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One defines the *density of states* measure, $dk(E)$ (and *integrated density of states*, aka IDS, $k(E) = dk((-\infty, E))$), by

$$dk = \frac{1}{p} \sum_{v \in V} d\mu_v$$

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$$dk = \frac{1}{p} \sum_{v \in V} d\mu_v$$

The big theorem of Sunada (1992), called *gap labelling*, says the following



The DOS and Gap Labelling

Theorem [Sunada] In any gap of the spectrum of $H_{\mathcal{T}}$, the IDS is an integral multiple of $1/p$. In particular, the spectrum has at most p connected components.

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The last statement has a simple proof given the first sentence.

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The last statement has a simple proof given the first sentence. Because dk is a finite sum, one sees that $\text{spec}(H_{\mathcal{T}}) = \text{supp}(dk)$.

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Theorem [Sunada] In any gap of the spectrum of $H_{\mathcal{T}}$, the IDS is an integral multiple of $1/p$. In particular, the spectrum has at most p connected components.

The last statement has a simple proof given the first sentence. Because dk is a finite sum, one sees that $\text{spec}(H_{\mathcal{T}}) = \text{supp}(dk)$. Thus if $a, b \notin \text{spec}(H_{\mathcal{T}})$ with spectrum in between, we must have that $k(b) > k(a)$ while, of course, k is constant in each gap.

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The Sunada Proof

I'll end this introduction by saying a little about Sunada's proof. It uses a deep theorem of Pimsner-Voiculescu (1982).

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I'll end this introduction by saying a little about Sunada's proof. It uses a deep theorem of Pimsner-Voiculescu (1982). Consider the homogeneous tree of degree 2ℓ which is the Cayley graph of \mathbb{F}_ℓ .

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Sunada first noted that in the general context of operators on trees of the type we looked at, the sum of diagonal matrix elements, one from each equivalence class of vertices, for operators commuting with our action of \mathbb{F}_ℓ (a natural von Neumann algebra) defines a natural normalized trace.



The Sunada Proof

$k(E)$ is just this normalized trace applied to the spectral projection $P_{(-\infty, E)}(H_{\mathcal{T}})$. This projection lies in the C^* -algebra generated by $H_{\mathcal{T}}$ if the projection is of the form $f(H_{\mathcal{T}})$ for a continuous function f and this is true if E is in a gap.

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The Sunada Proof

For period p , Sunada showed these C^* - and von Neumann algebras were twisted tensor products of the $p \times p$ matrices and the $p = 1$ -algebra.

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The Sunada Proof

For period p , Sunada showed these C^* - and von Neumann algebras were twisted tensor products of the $p \times p$ matrices and the $p = 1$ -algebra. Since the normalized trace of projections in $p \times p$ matrices is $1/p$, he could prove the normalized trace of any projections in the twisted tensor product had the same property which gives his theorem as a consequence of Pimsner-Voiculescu.

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The Pimsner-Voiculescu theorem is proven by them by using an exact sequence of K -theory groups.

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The Pimsner-Voiculescu theorem is proven by them by using an exact sequence of K -theory groups. While Effros and others have a simpler proof of their theorem, there is no elementary proof. This ends the introduction. My goal in the rest of the talk is our new proof of Sunada's gap labelling theorem which is so elementary we think of it as "the proof from the book"

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Hill's Equation

Our next subject, also preliminary, is on older results called gap labelling (all related to what we are calling gap labelling) focusing on the special case of the finite graph which is a simple cycle

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Our next subject, also preliminary, is on older results called gap labelling (all related to what we are calling gap labelling) focusing on the special case of the finite graph which is a simple cycle so the tree is just \mathbb{Z} and the Jacobi parameters are ordinary two-sided periodic sequences.

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$$-u''(x) + V(x)u(x) = \lambda u(x)$$

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Hill's equation is the differential equation on \mathbb{R}

$$-u''(x) + V(x)u(x) = \lambda u(x)$$

where λ is a (usually) real parameter and V is a real periodic function, i.e. $V(x + L) = V(x)$ for all real x and some $L \geq 0$.

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Hill's Equation

This is, of course, the $1D$ Schrödinger equation but Hill (the rare great 19^{th} century American scientist (he worked at Rutgers, the NJ state university, whose math building is named after him) was 40 years before Schrödinger and was studying perturbations of the moon orbit.

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$$-a_{n+1}u(n+1) + b_nu(n) + a_nu(n-1) = \lambda u(n)$$

which is of course a periodic Jacobi matrix on a tree where the tree is \mathbb{Z} (and if the period is p , the finite graph is a cyclic of length p .)

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We remark there is a version of Floquet theory and gap labelling for such periodic operators on \mathbb{Z}^ν - an abelian extension of the $1D$ theory as opposed to the tree theory which is a non-Abelian extension.



Floquet solution and Floquet multipliers

Solutions of the second order difference equation exist and are unique given $(u(0), u(1))$.

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Solutions of the second order difference equation exist and are unique given $(u(0), u(1))$. One can write the solution in terms of a 2×2 matrix

$$T(m; \lambda) \begin{pmatrix} u(1) \\ a(0)u(0) \end{pmatrix} = \begin{pmatrix} u(m+1) \\ a(m)u(m) \end{pmatrix}$$

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where $T(kp; \lambda) = T(p; \lambda)^k$ because of periodicity. By a simple calculation, $\det(T(p; \lambda)) = 1$, so the two eigenvalues of $T(p; \lambda)$,

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where $T(kp; \lambda) = T(p; \lambda)^k$ because of periodicity. By a simple calculation, $\det(T(p; \lambda)) = 1$, so the two eigenvalues of $T(p; \lambda)$, called *Floquet eigenvalues* denoted $\alpha_{\pm}(\lambda)$, obey $\alpha_{-}(\lambda) = \alpha_{+}(\lambda)^{-1}$.

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The solutions of the difference equation on \mathbb{Z} , u_{\pm} , with initial conditions the eigenfunctions of $T(p; \lambda)$ are called *Floquet solutions*.

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Floquet solution and Floquet multipliers

One subtlety we've avoided is that when $\lambda_{\pm} = \pm 1$, the two "eigenvalues" are equal so we can have geometric multiplicity 1 or 2.

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Floquet solution and Floquet multipliers

One subtlety we've avoided is that when $\lambda_{\pm} = \pm 1$, the two "eigenvalues" are equal so we can have geometric multiplicity 1 or 2. That is, for other values of λ_{\pm} , there are two Floquet solutions but for this case there might be either 1 or 2.

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Moreover, it is known that $\text{spec}(H)$ is precisely the set of λ for which there is a polynomially bounded solution, i.e. points where the Floquet eigenvalue has magnitude 1 rather than points where $\alpha_{+}(\lambda) > 1$.

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Moreover, it is known that $\text{spec}(H)$ is precisely the set of λ for which there is a polynomially bounded solution, i.e. points where the Floquet eigenvalue has magnitude 1 rather than points where $\alpha_{+}(\lambda) > 1$. The regions where there are bounded solutions are called *regions of stability* and where there aren't are *regions of instability*.

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Almost Periodic Gap Labelling

Before leaving the $1D$ case, we should mention a result going back to the 1980's that popularized the name "gap labelling" in a related but distinct context, namely for almost periodic classical Jacobi matrices, where a_n and b_n are almost periodic rather than periodic.

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$a_n \equiv 1, b_n = \beta \cos(\pi \alpha n + \theta)$ for parameters β, α, θ with α irrational. For this model, gap labelling says that in a gap, $k(\lambda) = m\alpha + n$ for integers m and n (and for the general case, it lies in the frequency module of the a, b .)

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The point is that the set of possible values is dense in $[0, 1]$ so that if all (or many) values occur, the spectrum is a Cantor set.

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The point is that the set of possible values is dense in $[0, 1]$ so that if all (or many) values occur, the spectrum is a Cantor set. The famous *ten martini problem* (which is a theorem of Avila-Jitomirskaya with important partial results by others, especially Puig) is that for all $\beta \neq 0$ and all irrational α , the almost Mathieu spectrum is a Cantor set.



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That concludes the background and we turn to our new proof.

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That concludes the background and we turn to our new proof. We recall that *Green's function* is the name given to matrix elements of the resolvent of a basic operator. Specifically, in our case, given $e \in \mathbb{E}(\Gamma)$, we pick any $\tilde{e} \in \mathbb{E}(\mathcal{T})$ with $\pi(\tilde{e}) = e$ and let $G_e(z)$ be the matrix element of $(H_{\mathcal{T}} - z)^{-1}$ corresponding to the two vertices of any \tilde{e} with $\pi(\tilde{e}) = e$.

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Dropping the edge \tilde{e} from \mathcal{T} breaks $\ell^2(\mathcal{T}) = \ell^2(\mathcal{T}_{\tilde{e}^-}) \oplus \ell^2(\mathcal{T}_{\tilde{e}^+})$ where $\ell^2(\mathcal{T}_{\tilde{e}^+})$ is the subspace with $\tau(\tilde{e})$ and $\ell^2(\mathcal{T}_{\tilde{e}^-})$ is the subspace with $\sigma(\tilde{e})$.

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$$m_e(z) = \langle \delta_{\tau(\tilde{e})}, (H_{\tilde{e}}^+ - z)^{-1} \delta_{\tau(\tilde{e})} \rangle$$



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Dropping the edge \tilde{e} from \mathcal{T} breaks $\ell^2(\mathcal{T}) = \ell^2(\mathcal{T}_{\tilde{e}^-}) \oplus \ell^2(\mathcal{T}_{\tilde{e}^+})$ where $\ell^2(\mathcal{T}_{\tilde{e}^+})$ is the subspace with $\tau(\tilde{e})$ and $\ell^2(\mathcal{T}_{\tilde{e}^-})$ is the subspace with $\sigma(\tilde{e})$. We let $H_{\tilde{e}}^{\pm}$ be the operators on $\ell^2(V(\mathcal{T}_{\tilde{e}}^{\pm}))$ with the restricted Jacobi parameters and set

$$m_e(z) = \langle \delta_{\tau(\tilde{e})}, (H_{\tilde{e}}^+ - z)^{-1} \delta_{\tau(\tilde{e})} \rangle$$

and, of course, $m_{\tilde{e}}(z) = \langle \delta_{\sigma(\tilde{e})}, (H_{\tilde{e}}^- - z)^{-1} \delta_{\sigma(\tilde{e})} \rangle$.



Schur Complements

The relation between resolvents of direct sums and resolvents of the summands was studied by Schur (1917) and is named the theory of Schur complements (called the method of Feshbach (1962) projections by theoretical physicists!).

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Schur Complements

The relation between resolvents of direct sums and resolvents of the summands was studied by Schur (1917) and is named the theory of Schur complements (called the method of Feshbach (1962) projections by theoretical physicists!). Applied to m - and Green's functions, this gives

$$\frac{1}{G_u(z)} = -z + b_u - \sum_{f \in \tilde{E}: \sigma(f)=u} a_f^2 m_f(z)$$

$$\frac{1}{m_f(z)} = -z + b_u - \sum_{\substack{f' \in \tilde{E}, f' \neq f \\ \sigma(f')=\tau(f)}} a_{f'}^2 m_{f'}(z)$$

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Schur Complements

which implies for any $e \in \tilde{E}$ that

$$G_{\sigma(e)} = \frac{1}{m_{\tilde{e}}^{-1} - a_e^2 m_e} = \frac{m_{\tilde{e}}}{1 - a_e^2 m_e m_{\tilde{e}}}$$

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This suggests a useful object

$$Q_e(z) = \frac{1}{1 - a_e^2 m_e(z) m_{\tilde{e}}(z)} = \frac{G_{\sigma(e)}(z)}{m_{\tilde{e}}(z)} = \frac{G_{\tau(e)}(z)}{m_e(z)}$$

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These equations imply the important result (of Chomsky-Schützenberger (1963) - that Chomsky!!),

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These equations imply the important result (of Chomsky-Schützenberger (1963) - that Chomsky!!), that G and m are algebraic functions, which implies (Avni-Breuer-Simon) that there is no singular continuous spectrum.

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The Floquet Function

The key to the new proof is a remarkable equality involving a new function we introduced and called the *Floquet function* defined by

$$\Phi(z) = \exp \left(p \int \log(t - z) dk(t) \right)$$

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defined originally in the upper half z -plane which clearly has an analytic continuation to a neighborhood of $\mathbb{C}_+ \cup (\mathbb{R} \setminus \text{spec}(H))$.

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Careful analysis of the imaginary part of the log in the gap, shows that the imaginary part of this last integral if $-p\pi k(E_0)$ is E_0 is a point in a gap.

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A Remarkable Formula

The key to our proof is the following formula involving Φ , G and m (or and Q)

$$\Phi(z) = \frac{\prod_{e \in E(\mathcal{G})} Q_e(z)}{\prod_{u \in V(\mathcal{G})} G_u(z)}$$

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where the second equality is just the definition of Q (and the first will be discussed shortly).

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Proof of the Magic Formula

Call the right side $\Psi(x)$.

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Proof of the Magic Formula

Call the right side $\Psi(x)$. The first thing to prove is that as $x \rightarrow \infty$ in $(0, \infty)$, one has that, $\Psi(-x) = x^p + O(x^{p-1})$ and the same for $\Phi(-x)$.

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Using the Schur complement formula for G_u^{-1} , one gets a formula for $(\log(G_u))'$ and from the definition of Q_e a formula for $(\log(Q_e))'$ from which one sees that $\sum_{e \in E} (\log(Q_e))' = \sum_{u \in V} [-G_u + (\log(G_u))']$ so that

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$$\sum_{e \in E} (\log(Q_e))' - \sum_{u \in V} (\log(G_u))' = \sum_{u \in V} -G_u$$

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Call the right side $\Psi(x)$. The first thing to prove is that as $x \rightarrow \infty$ in $(0, \infty)$, one has that, $\Psi(-x) = x^p + O(x^{p-1})$ and the same for $\Phi(-x)$. It follows that in that regime, $\log(\Phi/\Psi) = O(1/x)$, so it suffices to prove that $\log(\Phi)$ and $\log(\Psi)$ have the same derivative.

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The left side is just $[\log(\Psi)]'$ and the right is $[\log(\Phi)]'$ proving the magic formula.

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Proof of Gap Labelling

It is easy to see that the operators H_e^\pm have essential spectra which are subsets of the essential spectra of H_T

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Proof of Gap Labelling

It is easy to see that the operators H_e^\pm have essential spectra which are subsets of the essential spectra of H_T (in fact, one can show the essential spectra are equal).

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Proof of Gap Labelling

It is easy to see that the operators H_e^\pm have essential spectra which are subsets of the essential spectra of H_T (in fact, one can show the essential spectra are equal). It follows that these operators have only discrete spectrum in gaps of H_T so the Green's and m -functions only have isolated zeros and poles in the gaps

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Proof of Gap Labelling

It is easy to see that the operators H_e^\pm have essential spectra which are subsets of the essential spectra of H_T (in fact, one can show the essential spectra are equal). It follows that these operators have only discrete spectrum in gaps of H_T so the Green's and m -functions only have isolated zeros and poles in the gaps (indeed, using that they are algebraic, only finitely many). Thus the function, Ψ in the magic formula, which is built from G_u 's and m_e 's, is regular, real, and non-zero at most points in a gap.

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Proof of Gap Labelling

It is easy to see that the operators H_e^\pm have essential spectra which are subsets of the essential spectra of H_T (in fact, one can show the essential spectra are equal). It follows that these operators have only discrete spectrum in gaps of H_T so the Green's and m -functions only have isolated zeros and poles in the gaps (indeed, using that they are algebraic, only finitely many). Thus the function, Ψ in the magic formula, which is built from G_u 's and m_e 's, is regular, real, and non-zero at most points in a gap. Thus, by the magic formula, the argument of Φ , which we have seen is $-p\pi k(E)$ (for E in the gap) is an integral multiple of π which is gap labelling!!

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Point Spectrum

If there is time, I'll say a few word about our other new proof.

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Point Spectrum

If there is time, I'll say a few word about our other new proof. In the $1D$ case, H does not have any point spectrum, but in other cases that is not true for example, one Γ of Aomoto where $r < g$ are fixed positive integers.

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Statement

Given an eigenvalue, λ , define $X_1(\lambda)$ to be the set of vertices, $v \in V$, so that for some \tilde{v} with $\pi(\tilde{v}) = v$ there is some eigenfunction ψ associated to λ , with $\psi(\tilde{v}) \neq 0$. Define $\partial X_1(\lambda)$ to be those $v \in V$ not in $X_1(\lambda)$ but neighbors of points in $X_1(\lambda)$, and we let $E(\lambda)$ be the set of edges with both endpoints in $X_1(\lambda)$.

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Theorem [Aomoto Index Theorem] The measure dk has a mass at an eigenvalue, λ , of weight $I(\lambda)/p$ where

$$I(\lambda) = \#(X_1(\lambda)) - \#(\partial X_1(\lambda)) - \#(E(\lambda))$$

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Theorem [Aomoto Index Theorem] The measure dk has a mass at an eigenvalue, λ , of weight $I(\lambda)/p$ where

$$I(\lambda) = \#(X_1(\lambda)) - \#(\partial X_1(\lambda)) - \#(E(\lambda))$$

A consequence of this theorem is that if Γ has a fixed degree (equivalently, \mathcal{T} does), then $H_{\mathcal{T}}$ has no point spectrum.

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Remarks

The original proof of Aomoto is opaque and our new proof is simpler but there is a 2022 paper of Banks, Garza-Vargas and Mukhejee with a particularly elegant way to understand point spectrum including a lovely proof of the index theorem.

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Barry Simon

$xy \leq \frac{x^p}{p} + \frac{y^q}{q}$

$\hat{f}(\mathbf{k}) = (2\pi)^{-n/2} \int \exp(-i\mathbf{k} \cdot \mathbf{x}) f(\mathbf{x}) d^n x$

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 **Basic Complex Analysis**

Barry Simon

ANALYSIS
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2A


Simon

A Comprehensive Course in Analysis by Poincaré Prize winner Barry Simon is a five-volume set that can serve as a graduate-level analysis textbook with a lot of additional bonus information, including hundreds of problems and numerous notes that extend the text and provide important historical background. Depth and breadth of exposition make this set a valuable reference source for almost all areas of classical analysis.


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Part 2A is devoted to basic complex analysis. It interweaves three analytic threads associated with Cauchy, Riemann, and Weierstrass, respectively. Cauchy's view focuses on the differential and integral calculus of functions of a complex variable, with the key topics being the Cauchy integral formula and contour integration. For Riemann, the geometry of the complex plane is central, with key topics being fractional linear transformations and conformal mapping. For Weierstrass, the power series is king, with key topics being spaces of analytic functions, the product formulas of Weierstrass and Hadamard, and the Weierstrass theory of elliptic functions. Subjects in this volume that are often missing in other texts include the Cauchy integral theorem when the contour is the boundary of a Jordan region, continued fractions, two proofs of the big Picard theorem, the uniformization theorem, Ahlfors's function, the sheaf of analytic germs, and Jacobi, as well as Weierstrass, elliptic functions.

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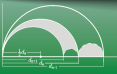
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Basic Complex Analysis
A Comprehensive Course in Analysis, Part 2A

Barry Simon



$$f(z_0) = \frac{1}{2\pi i} \int_{|z|=1} \frac{f(z)}{z - z_0} dz$$





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for Hill's
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Green's and
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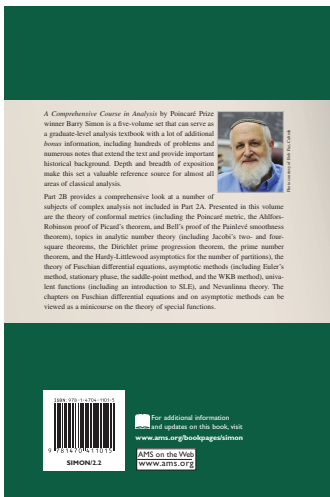


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Advanced Complex Analysis A Comprehensive Course in Analysis, Part 2B

Barry Simon

$$\frac{\pi(x)}{(x/\log x)} \rightarrow 1$$



$$J_n(x) \sim \sqrt{\frac{2}{\pi x}} \cos\left(x - \frac{\alpha x}{2} - \frac{\pi}{4}\right) + o(x^{-1/2})$$



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Part 3 returns to the themes of Part 1 by discussing pointwise limits (going beyond the usual focus on the Hardy-Littlewood maximal function by including ergodic theorems and martingale convergence), harmonic functions and potential theory, frames and wavelets, H^p spaces (including bounded mean oscillation (BMO)) and, in the final chapter, lots of inequalities, including Sobolev spaces, Calderón-Zygmund estimates, and hypercontractive semigroups.



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Harmonic Analysis

ANALYSIS

Part
3

Simon

Harmonic Analysis
A Comprehensive Course in Analysis, Part 3

Barry Simon



$$\|f - f_Q\|_Q = \frac{1}{|Q|} \int_Q |f(x) - f_Q| dx$$

$$|\{x \mid M_{\text{HL}} f(x) > \alpha\}| \leq \frac{3^n}{\alpha} \|f\|_{L^1(\mathbb{R}^n, dx)}$$



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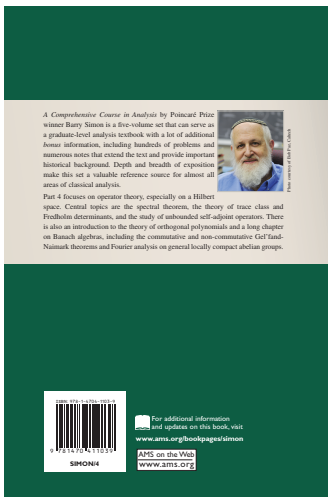
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$$A = \int t dE_t$$

$$\det(1 + zA) = \prod_{k=1}^{N(A)} (1 + z\lambda_k(A))$$

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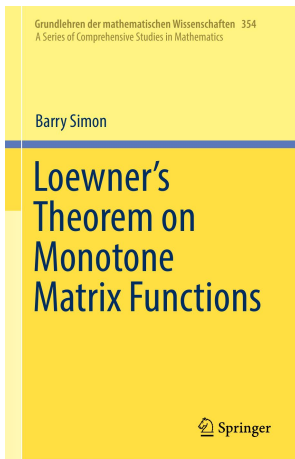
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