

**OPERATORS WITH SINGULAR CONTINUOUS SPECTRUM:  
III. ALMOST PERIODIC SCHRÖDINGER OPERATORS**

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ABSTRACT. We prove that one-dimensional Schrödinger operators with even almost periodic potential have no point spectrum for a dense  $G_\delta$  in the hull. This implies purely singular continuous spectrum for the almost Mathieu equation for coupling larger than 2 and a dense  $G_\delta$  in  $\theta$  even if the frequency is an irrational with good Diophantine properties.

**§1. Introduction**

This is a paper that provides yet another place where singular continuous spectrum occurs in the theory of Schrödinger operators and Jacobi matrices (see [5,6,2,10,3]). It is especially interesting because it will provide examples where a non-resonance condition in a KAM argument is not merely needed for technical reasons but necessary.

Our main results, proven in §2, do not deal directly with singular continuous spectrum but only with continuous spectrum.

**Theorem 1S.** *Let  $V$  be an even almost periodic function on  $(-\infty, \infty)$  and let  $\Omega$  be the hull of  $V$  and  $V_\omega(x)$  the corresponding function for  $\omega \in \Omega$ . Then there is a dense  $G_\delta, U$  in  $\Omega$  (in the natural metric topology), so that if  $\omega \in U$ , then  $H_\omega \equiv \frac{d^2}{dx^2} + V_\omega(x)$  has no eigenvalues as an operator on  $L^2(\mathbb{R})$ .*

For the Jacobi case, we let  $h_0$  be the operator on  $\ell^2(\mathbb{Z})$  defined by  $(h_0u)(n) = u(n+1) + u(n-1)$ .

**Theorem 1J.** *Let  $V$  be an even almost periodic function on  $\mathbb{Z}$ ,  $\Omega$  its hull, and  $V_\omega(n)$  the function associated to  $\omega \in \Omega$ . Then there is a dense  $G_\delta, U$  in  $\Omega$  so that if  $\omega \in U$ , then  $H_\omega = h_0 + V_\omega(n)$  has no eigenvalues as an operator on  $\ell^2(\mathbb{Z})$ .*

The  $G_\delta$  set  $U$  will be rather explicit—see §2. By combining this with the machinery of [10], we can sometimes get singular continuous spectrum.

**Theorem 2.** *In the context of Thm. 1, suppose there is a single  $\omega \in \Omega$  so that  $H_\omega$  has no absolutely continuous spectrum. Then for a dense  $G_\delta, \tilde{U}$ ,  $H_\omega$  has purely singular continuous spectrum.*

*Proof.* Let  $U_1 = \{\omega \in \Omega \mid H_\omega \text{ has no a.c. spectrum}\}$ . By [10],  $U_1$  is a  $G_\delta$ . By hypothesis,  $\omega_0$  and its translates lie in  $U_1$ , so  $U_1$  is a dense  $G_\delta$ . Thus,  $\tilde{U} = U_1 \cap U$  is a dense  $G_\delta$ .

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**Example 1.** Consider the Jacobi matrix with

$$V_\theta(n) = \lambda \cos(\pi\beta n + \theta). \quad (1)$$

If  $\lambda > 2$ , the Lyapunov exponent is positive ([1,7]) so if  $\beta$  is irrational, there is no a.c. spectrum for Lebesgue a.e.  $\theta$  (see e.g. [1]), so  $h_\theta$  has purely singular continuous spectrum for a dense  $G_\delta$  of  $\theta$ .

Sinai [11] and Fröhlich-Spencer-Wittwer [4] have proven for  $\lambda$  large and  $\beta$  having good Diophantine properties, a.e.  $\theta$  has pure point spectrum, and Jitomirskaya [8] has proven that for  $\lambda \geq 15$ . In that case there are intertwined locally uncountable sets of  $\theta$  with only pure point and with only singular continuous spectrum. For  $\lambda = 2$ ,  $\text{spec}(h_\theta)$  has zero measure for many irrational  $\beta$ 's [9] and so no a.c. spectrum. We conclude

**Theorem 3.** *For the example (1),  $h_\theta$  has purely singular continuous spectrum for a dense  $G_\delta$  of  $\theta$ 's if  $\beta$  is irrational and  $\lambda > 2$  or if the continued fraction expansion of  $\beta$  has unbounded integers and  $\lambda = 2$ .*

**Example 2.** Consider the Schrödinger case with  $V_\theta(x) = -k[\cos(2\pi x) + \cos(2\pi\beta x + \theta)]$ . Then, Fröhlich-Spencer-Wittwer [4] have proven for a.e.  $\theta$  ( $k$  large enough), there is pure point spectrum for low energies. Sorets-Spencer [12] have proven positivity of the Lyapunov exponent for a wider area of low energy. We conclude that for a dense  $G_\delta$  of  $\theta$ , there is purely singular continuous spectrum for low energies.

## §2. Proof of Theorem 1

We'll consider the Jacobi case in detail and then discuss the changes for the Schrödinger case. Let  $V_{\omega_0}$  be the even almost periodic function on  $\mathbb{Z}$ :

$$V_{\omega_0}(-n) = V_{\omega_0}(n).$$

Fix once and for all a number  $B$  so

$$B > 4 \ln(3 + 2 \sup_n |V_{\omega_0}(n)|) \equiv 4 \ln \alpha. \quad (2.1)$$

$\alpha$  is chosen so that the matrix  $\begin{pmatrix} E - V(u) & 1 \\ 1 & 0 \end{pmatrix}$  has norm bounded by  $\alpha$  if  $|E| \leq 2 + \sup_n |V_{\omega_0}(n)|$ .

Let  $\Omega$  be the hull of  $V$ , that is, the closure in  $\|\cdot\|_\infty$  of translates of  $V$ ; it is compact by hypothesis. Define  $\rho$  on  $\Omega$  by

$$\rho(\omega, \omega') \equiv \sup_n (|V_\omega(n) - V_{\omega'}(n)|)$$

and define maps  $R$  and  $T$  on  $\Omega$  by

$$V_{R\omega}(n) = V_\omega(-n) \quad V_{T\omega}(n) = V_\omega(n-1).$$

**Lemma 2.1.** Let  $U_n = \bigcup_{|m|>n} \{\omega \mid \rho(RT^{2m}\omega, \omega) < e^{-B|m|}\}$  and let  $U = \bigcap_{n=1}^{\infty} U_n$ . Then  $U_n$  is a dense open set and  $U$  is a dense  $G_\delta$  in  $\Omega$ .

*Proof.* Let  $\omega_m = T^{-m}\omega_0$ . Then  $RT^{2m}\omega_m = \omega_m$  since  $R\omega_0 = \omega_0$ , so  $\omega_m \in U_n$  if  $|m| > n$ . It is easy to see the set of  $\{\omega_m \mid |m| > n\}$  is dense in  $\Omega$ , so  $U_n$  is dense. It is clearly open and so  $U = \bigcap U_n$  is a dense  $G_\delta$  by the Baire category theorem.

$U$  is the set of  $\omega$ 's for which there exists an infinite sequence  $m_i$  with  $|m_i| \rightarrow \infty$  with  $\rho(RT^{2m_i}\omega, \omega) < e^{-B|m_i|}$ . For a subsequence, either  $m_i \rightarrow \infty$  or  $m_i \rightarrow -\infty$  and by reflection invariance, we can suppose  $m_i \rightarrow \infty$ . Thus, Thm. 1J follows from

**Theorem 2.2.** Suppose that  $V$  is a function obeying

$$|V(2m_i - n) - V(n)| \leq e^{-Bm_i} \quad (2.2)$$

for a sequence  $m_i \rightarrow \infty$  where  $B$  is given by (2.1). Then

$$u(n+1) + u(n-1) + V(n)u(n) = Eu(n) \quad (2.3)$$

has no  $\ell^2$  solutions for any  $E$ .

*Remark.* The intuition behind the proof is that any  $u$  obeying (2.3) has to be close to being even or odd about  $m_i$  so  $u(n) \approx 0$ .

*Proof.* Suppose not. Then we can find a solution  $u$  of (2.3) in  $\ell^2$  which we normalize, so that

$$\sum_n |u(n)|^2 = 1. \quad (2.4)$$

We let  $u_i(n) \equiv u(2m_i - n)$ . Let  $W(f, g)(n) = f(n+1)g(n) - f(n)g(n+1)$  be the Wronskian as usual, and let

$$\Phi(n) = \begin{pmatrix} u(n+1) \\ u(n) \end{pmatrix}; \quad \Phi_i(n) = \begin{pmatrix} u_i(n+1) \\ u_i(n) \end{pmatrix}$$

as two component vectors.

**Step 1.** *Almost constancy of  $W(u, u_i)$*

By a standard calculation using (2.3)

$$\begin{aligned} |W(u, u_i)(n) - W(u, u_i)(n-1)| &\leq |V(n) - V(2m_i - n)| |u(n)u_i(n)| \\ &\leq e^{-Bm_i} \end{aligned} \quad (2.5)$$

by (2.2) and (2.4).

**Step 2.** *Smallness of  $W(u, u_i)$  for  $m_i$  large*

Since  $u$  and  $u_i$  are in  $\ell^2$  with  $\ell^2$  norm 1, the Schwarz inequality implies that  $\sum_n |W(n)| \leq 2$ . Thus for some  $n$  with  $|n| \leq e^{Bm_i/2}$ , we must have that  $|W(n)| \leq e^{-Bm_i/2}$ . By (2.5) we see that for  $|n| \leq e^{Bm_i/2}$ , we have that

$$|W(n)| \leq 3e^{-Bm_i/2} \quad (2.6)$$

and in particular for  $n = m_i$ .

Now define  $u_i^\pm = u \pm u_i$ ,  $\Phi_i^\pm = \Phi \pm \Phi_i$ .

**Step 3.** *Smallness of  $\Phi_i^+(m_i)$  or  $\Phi_i^-(m_i)$* 

Since  $W(u_i^-, u_i^+) = 2W(u, u_i)$  and  $u_i^-(m_i) = 0$ , we see that

$$|u_i^+(m_i)u_i^-(m_i + 1)| \leq 6e^{-Bm_i/2}$$

so either

$$|u^+(m_i)| \leq \sqrt{6}e^{-Bm_i/4} \quad (2.7)$$

or

$$|u_i^-(m_i + 1)| \leq \sqrt{6}e^{-Bm_i/4}. \quad (2.8)$$

We claim that this means either

$$\|\Phi_i^\pm(m_i)\| \leq Ce^{-Bm_i/4} \quad (\text{for one of } + \text{ or } -). \quad (2.9)$$

If (2.8) holds, (2.9) is immediate since  $u_i^-(m_i) = 0$ . If (2.7) holds, note that by (2.3)

$$u^+(m_i + 1) + \frac{1}{2}(V(m_i) - E)u^+(m_i) = 0$$

so (2.9) holds for  $\Phi_i^+$ .

**Step 4.** *Smallness of  $\Phi_i^\pm(0)$* 

Let  $T_i^{(1)}$  be the transfer matrix for (2.3), taking  $\Phi(m_i)$  to  $\Phi(0)$  and let  $T_i^{(2)}$  be the same with  $V(2m_i - n)$  so

$$\begin{aligned} T_i^{(1)}\Phi(m_i) &= \Phi(0) \\ T_i^{(2)}\Phi_i(m_i) &= \Phi_i(0). \end{aligned}$$

Writing out  $T_i$  as a product and using the definition of  $\alpha$  and (2.2), we have that

$$\|T_i^{(1)} - T_i^{(2)}\| \leq 2m_i\alpha^{m_i-1}e^{-Bm_i} \leq 2m_i e^{-3Bm_i/4}.$$

Writing

$$\begin{aligned} \Phi_i^\pm(0) &= T_i^{(1)}\Phi(m_i) \pm T_i^{(2)}\Phi_i(m_i) \\ &= T_i^{(1)}(\Phi_i^\pm(m_i)) \mp (T_i^{(1)} - T_i^{(2)})\Phi_i(m_i) \end{aligned}$$

we see that

$$\|\Phi_i^\pm(0)\| \leq m_i e^{-3Bm_i/4} + C(\alpha e^{-B/4})^{m_i}$$

goes to zero as  $m_i \rightarrow \infty$ .

**Step 5.** *Completion of the proof*

By the last fact,  $\|\Phi(0)\| - \|\Phi(2m_i)\| \rightarrow 0$  which is only consistent with  $u \in \ell^2$  if  $\|\Phi(0)\| = 0$  which implies that  $u = 0$ .

For the continuum (Schrödinger case), here are the changes: We can suppose (2.2) holds, but with  $e^{-Bm_i}$  replaced by  $e^{-m_i^2}$  (any  $f(m)$  with  $\lim_{i \rightarrow \infty} m_i^{-1} \ln f(m)^{-1} = \infty$  will do). We normalize  $u$  so that

$$\int [u(x)^2 + u'(x)^2] dx = 1. \quad (2.10)$$

**Step 1.** By (2.10) and a Sobolev estimate,  $u$  and  $u'$  are uniformly bounded so  $|\frac{dW}{dx}(u, u_i)(x)| \leq Ce^{-m_i^2}$  for some  $C$ .

**Step 2.**  $\int |W(u, u_i)| dx \leq 2$ , so, by the same argument

$$|W(u, u_i)|(x) \leq (2C + 1)e^{-m_i^2} \quad \text{if } |x| \leq e^{m_i^2/2}.$$

**Step 3.** This is actually easier since  $(u_i^+)'(m_i) = 0$  and  $u_i^-(m_i) = 0$ .

**Step 4.** This is similar. The transfer matrix is bounded by  $e^{Cm_i}$  where  $C$  is  $E$ -dependent (and goes to infinity as  $E \rightarrow \infty$ ) which is always beaten out by  $e^{-m_i^2/2}$ .

**Step 5** is unchanged.

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